The HI-to-H₂ Transition and HI columns in Galaxy Star-Forming Regions

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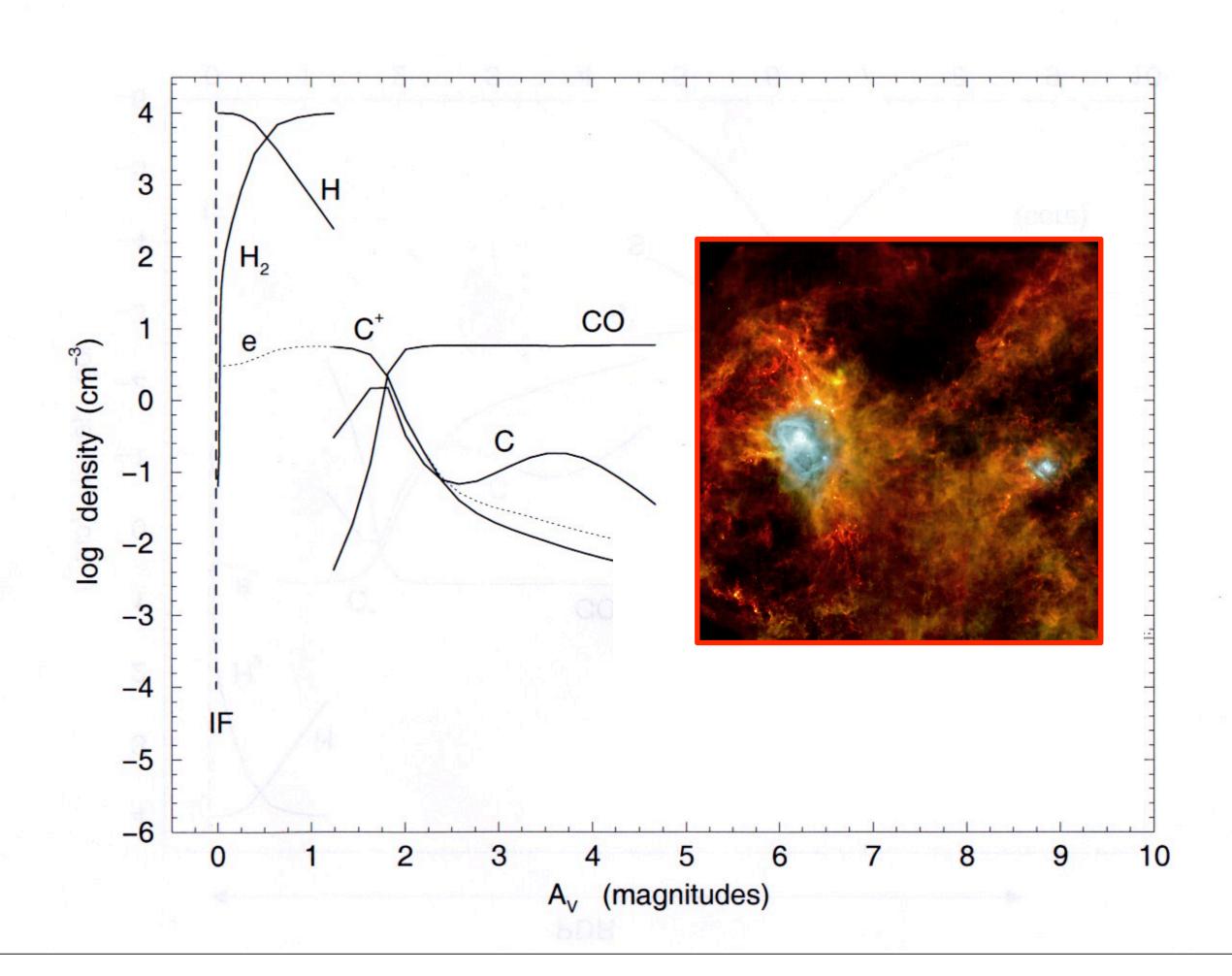
Star Formation Across Space and Time 14 November 2014

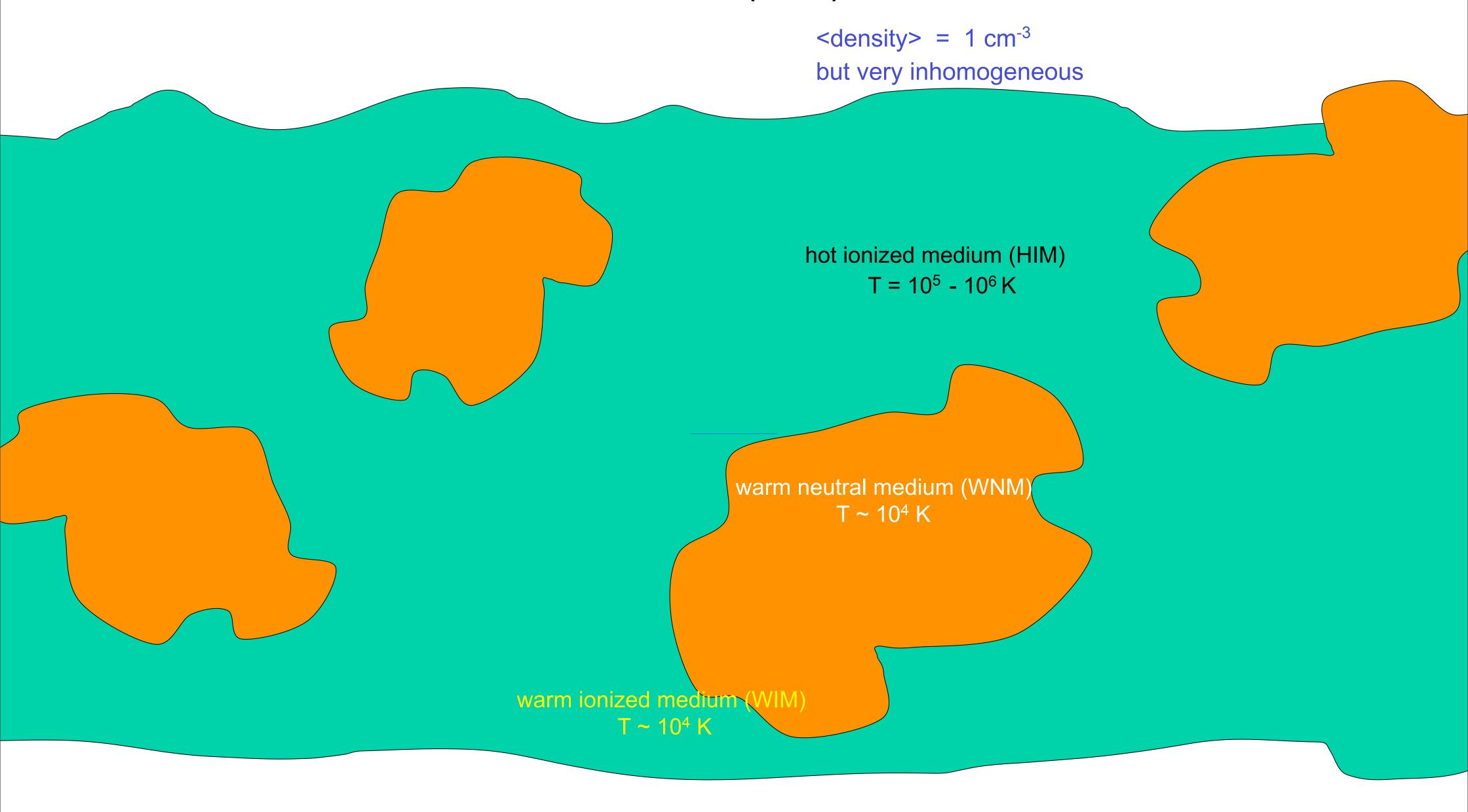


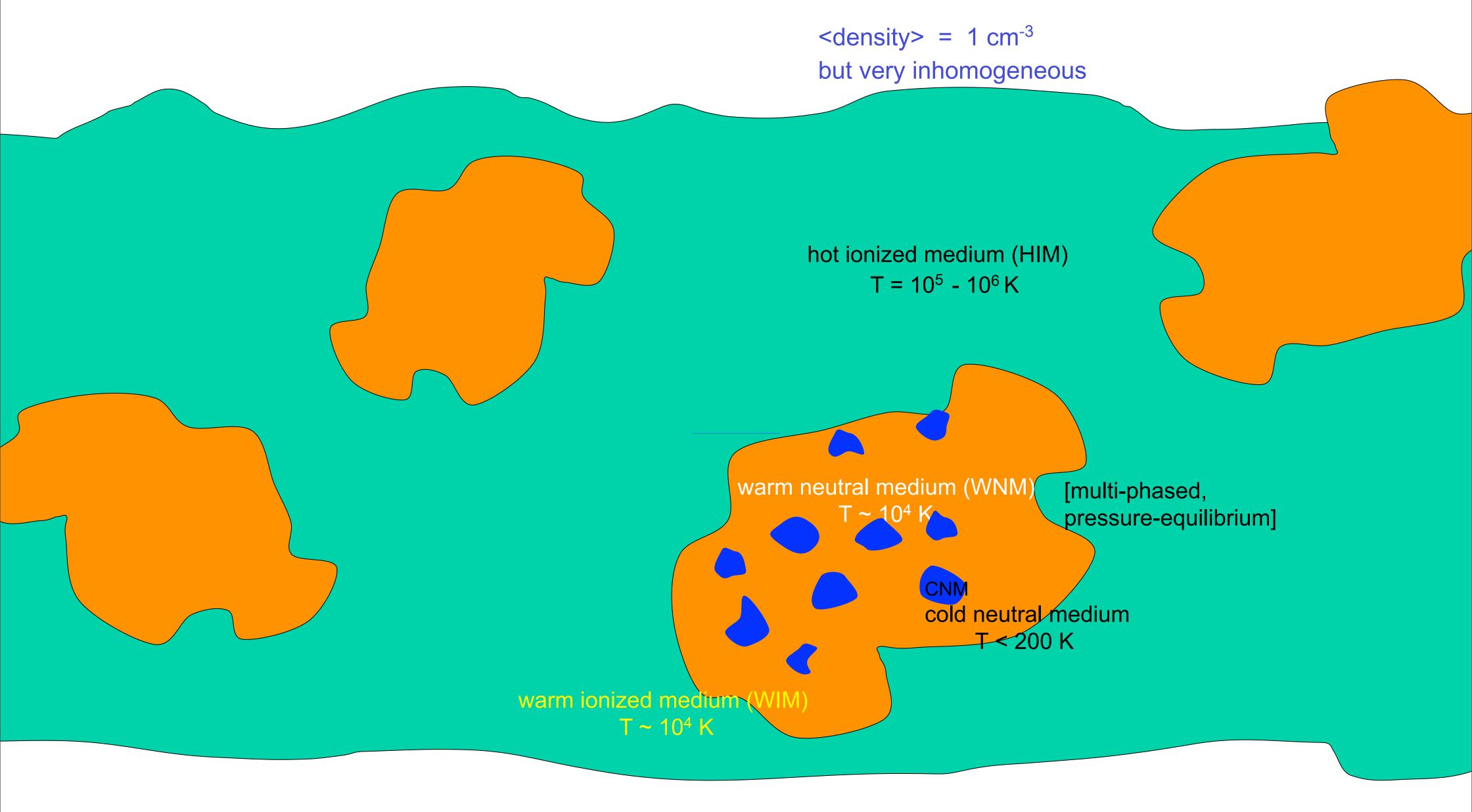
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Star Formation Across Space and Time 14 November 2014







Interstellar gas exposed to

starlight shock waves energetic particles (cosmic-rays) magnetic fields

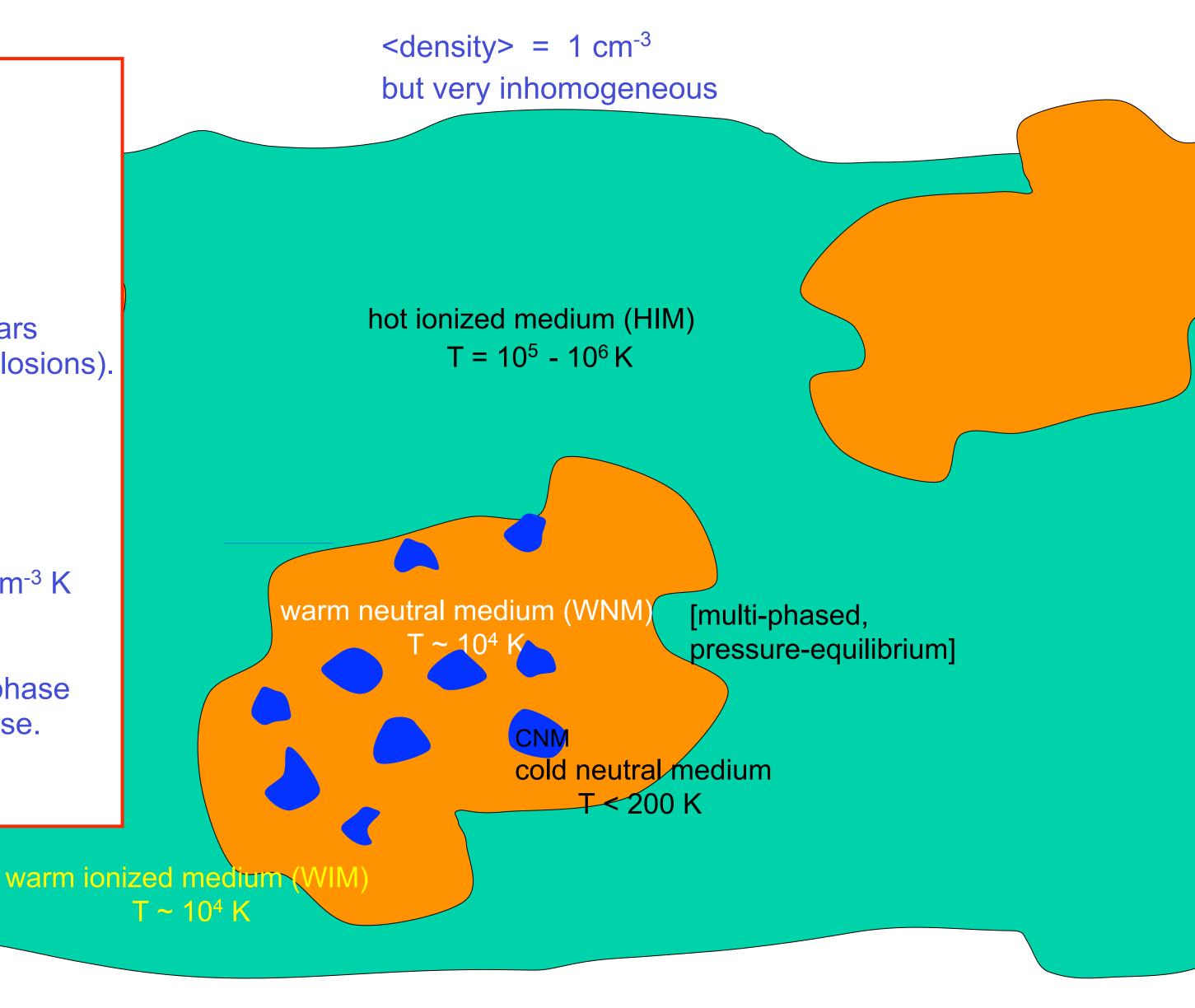
Global ISM heated and energized by stars (outflows, radiation, and supernova explosions).

Turbulent!

Total Galactic ISM mass = 5x10⁹ M_☉

Mid-plane thermal pressure = 2.5x10³ cm⁻³ K (at Solar circle)

95% of mass in cold neutral hydrogen phase 95% of volume in warm/hot ionized phase.



Interstellar gas exposed to

starlight shock waves energetic particles (cosmic-rays) magnetic fields

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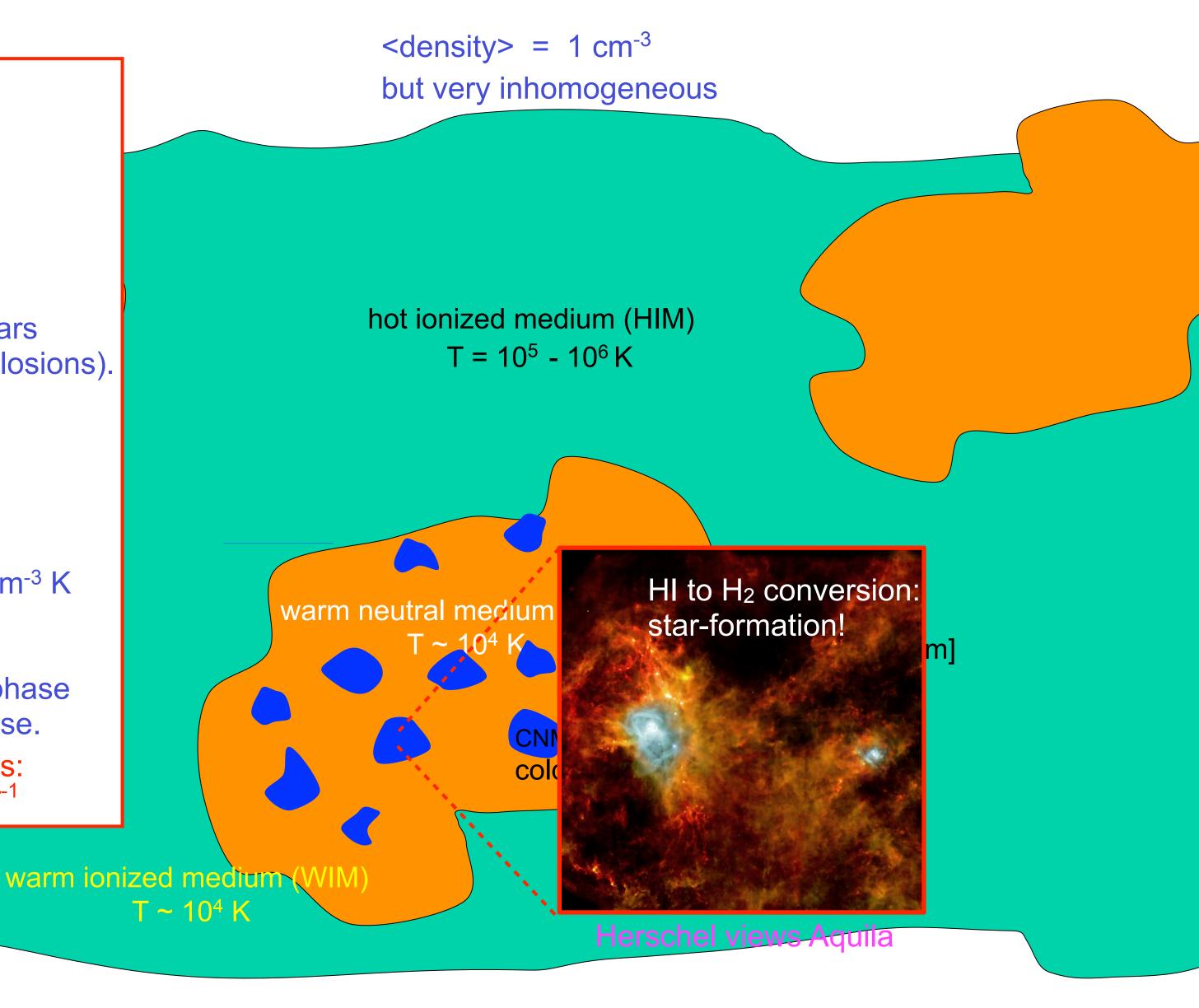
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95% of mass in cold neutral hydrogen phase 95% of volume in warm/hot ionized phase.

stars form in cold molecular (H₂) clouds: Galactic star-formation rate ~3 M_☉ yr⁻¹



Talk Outline:

- motivation.
- HI-to-H₂ transition, some radiative transfer computations.
- analytic formula for the HI column density.
- self-regulated media.
- observations: from Perseus to galaxies.

HI-to-H₂ Transitions and HI Column Densities in Galaxy Star-Forming Regions

Amiel Sternberg¹, Franck Le Petit², Evelyne Roueff,² and Jacques Le Bourlot²

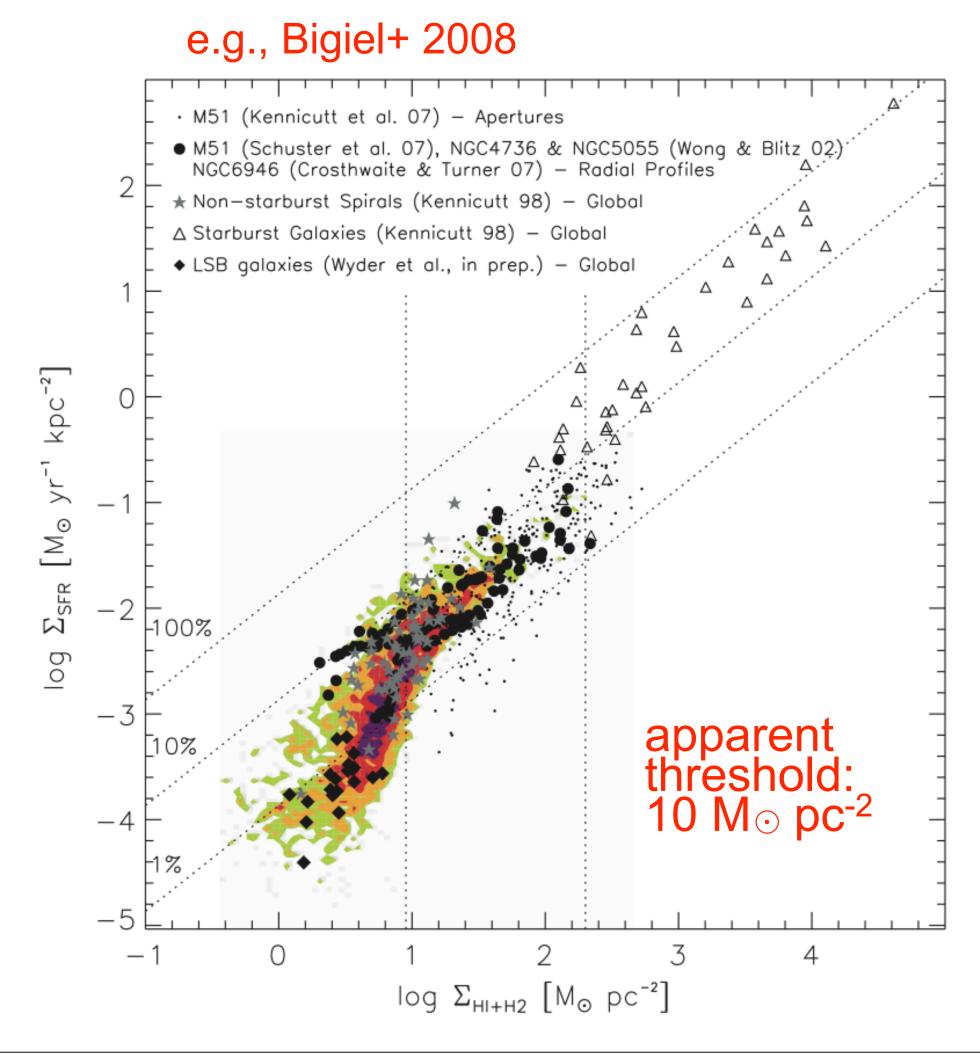


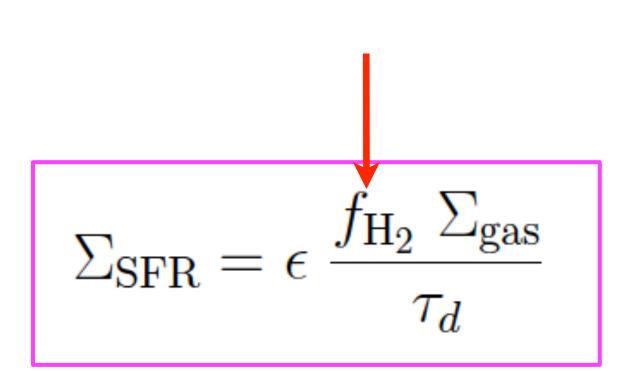
2014 ApJ 790 10



Kennicutt-Schmidt Relation:

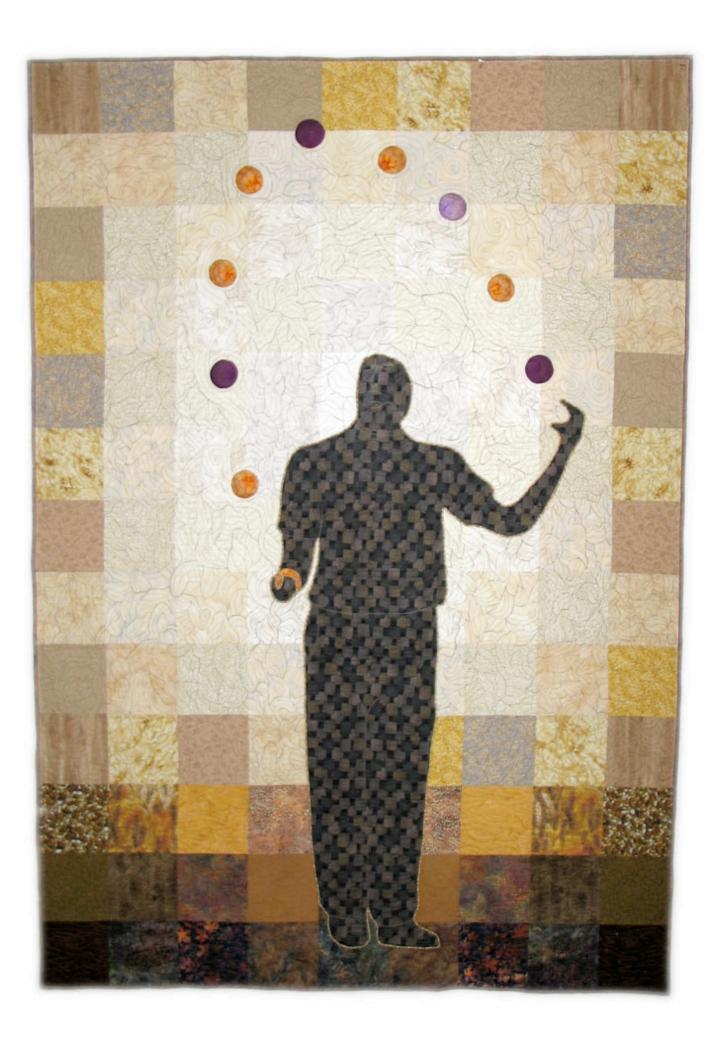
Schmidt 1959 ApJ 129 253 Kennicutt 1998 ApJ498 541 Genzel et al. 2010 MNRAS 407 2091 ["SINS(VLT)/IRAM" projects]





I am going to be juggling many parameters

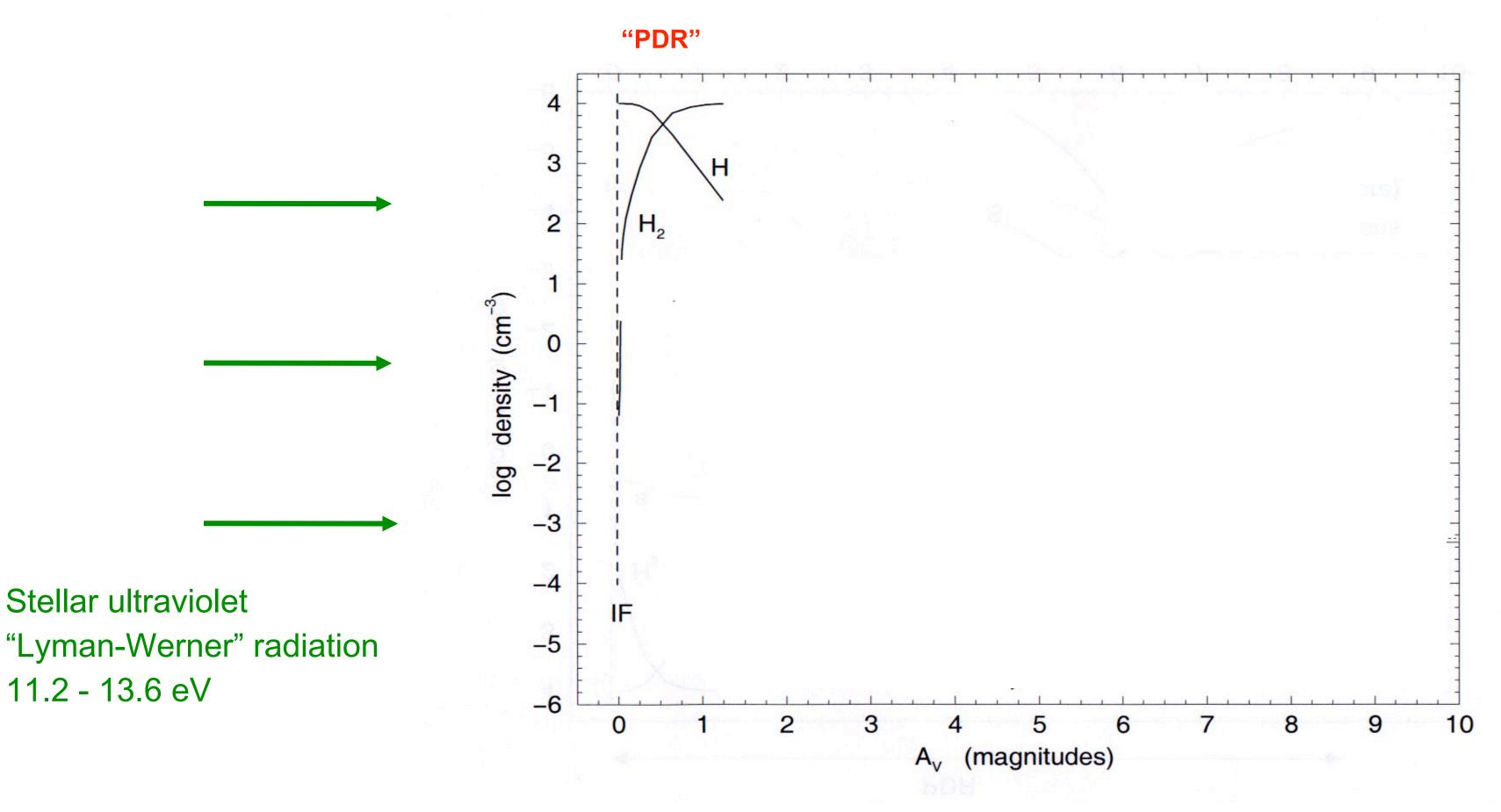
- keep your eye on them!



HI to H₂ Transition in Dense Star-Forming Molecular Clouds:

H₂ formation (by grain catalysis) versus far-UV photodissociation.

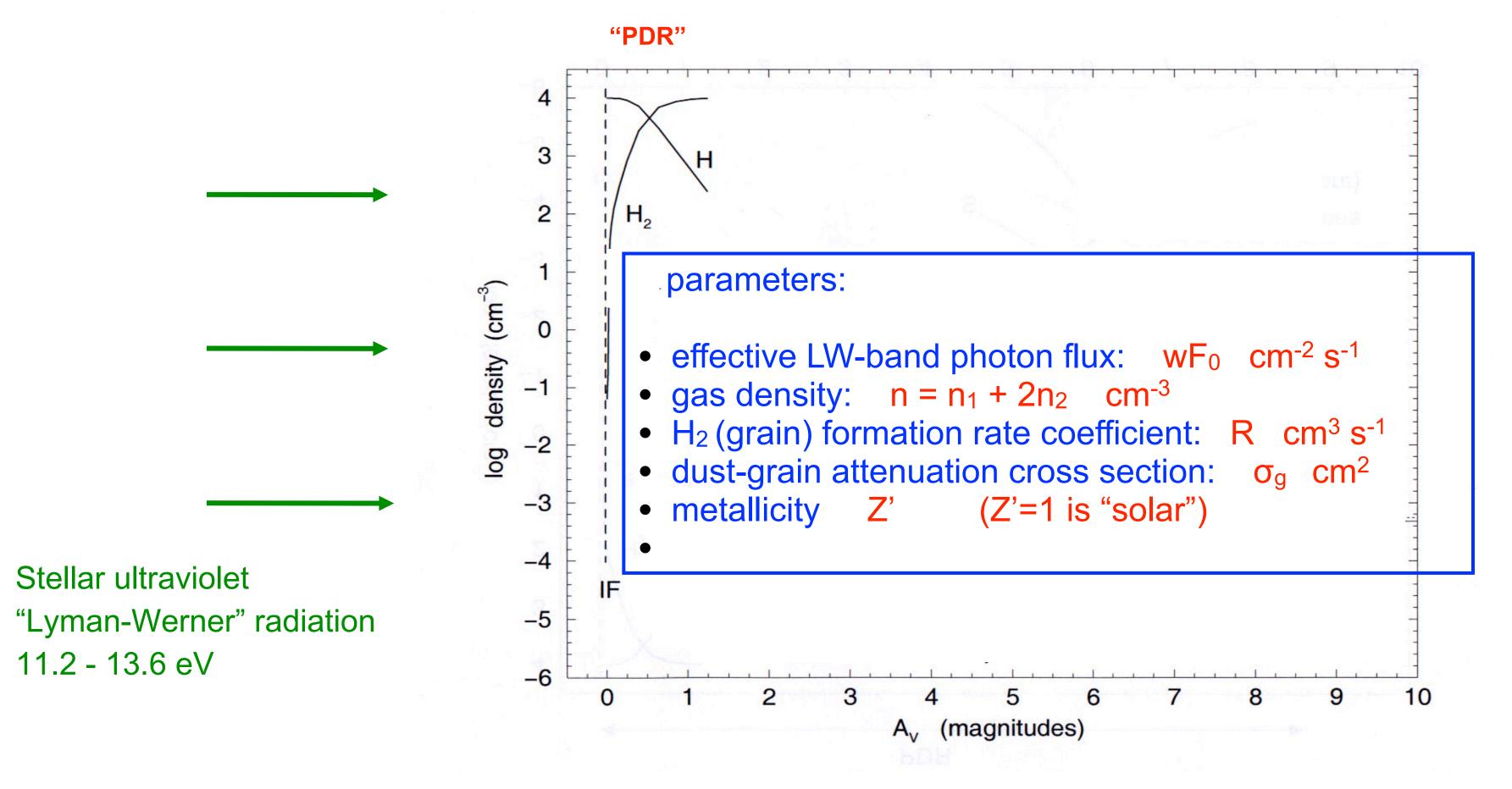
Shielding required.

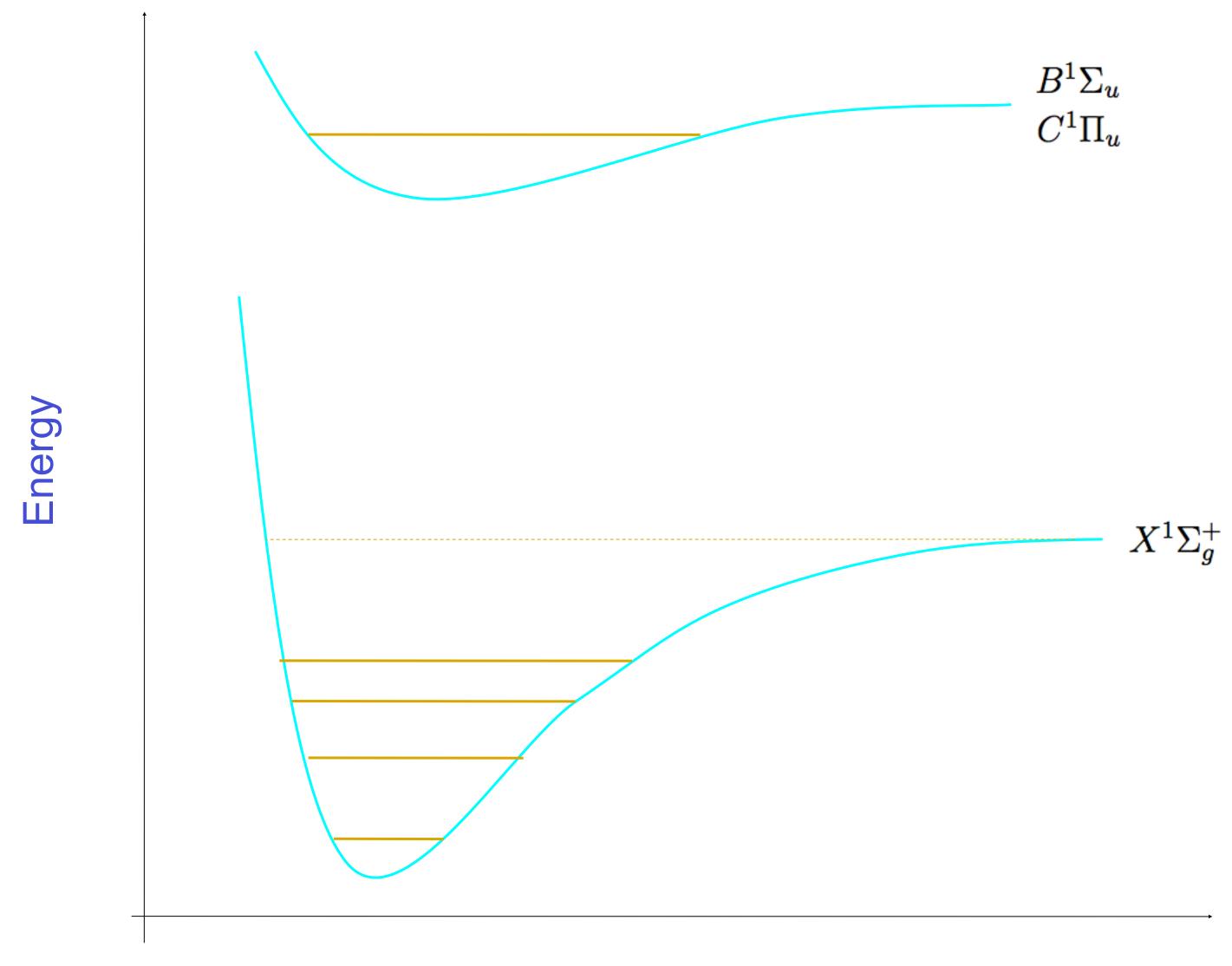


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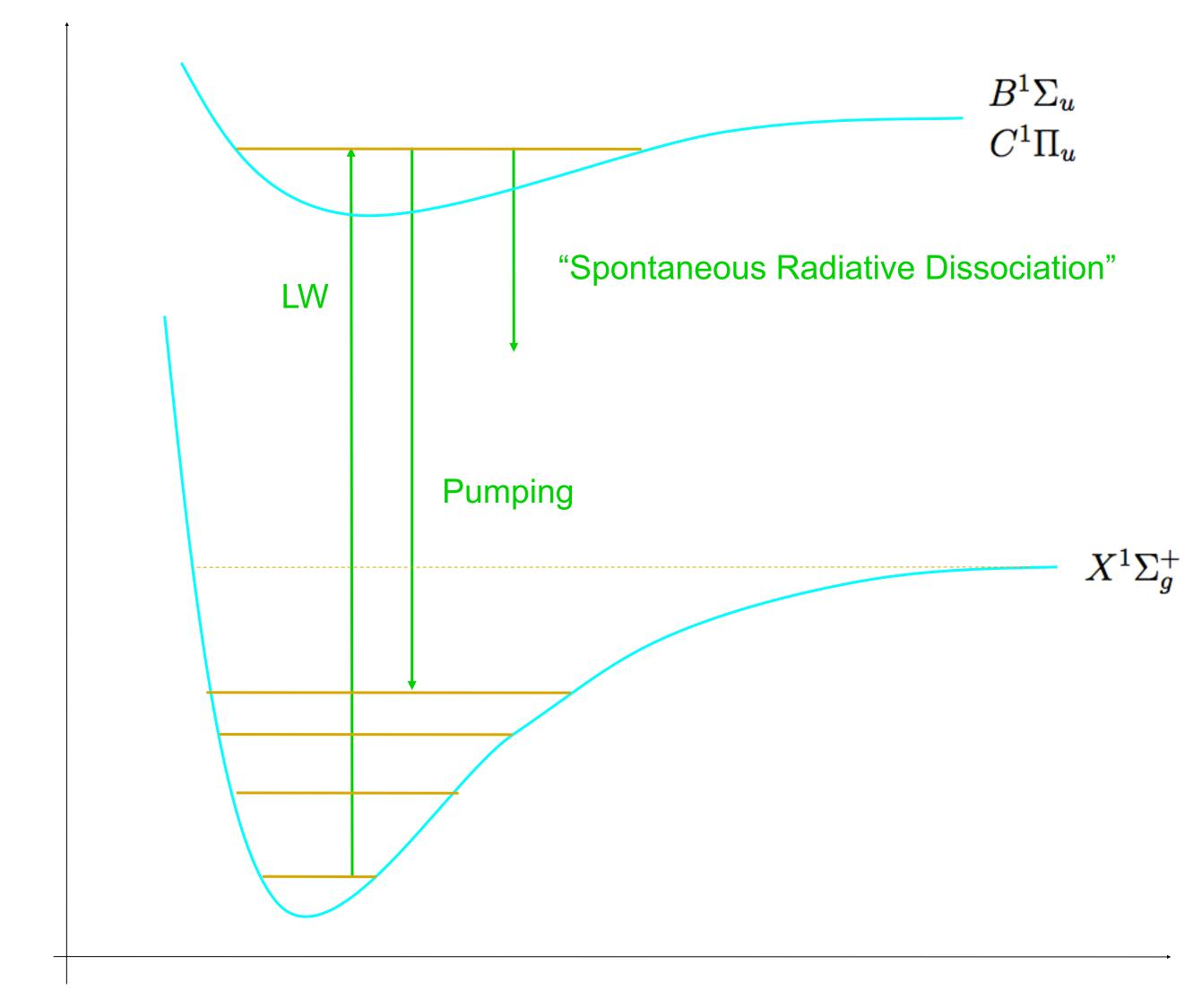
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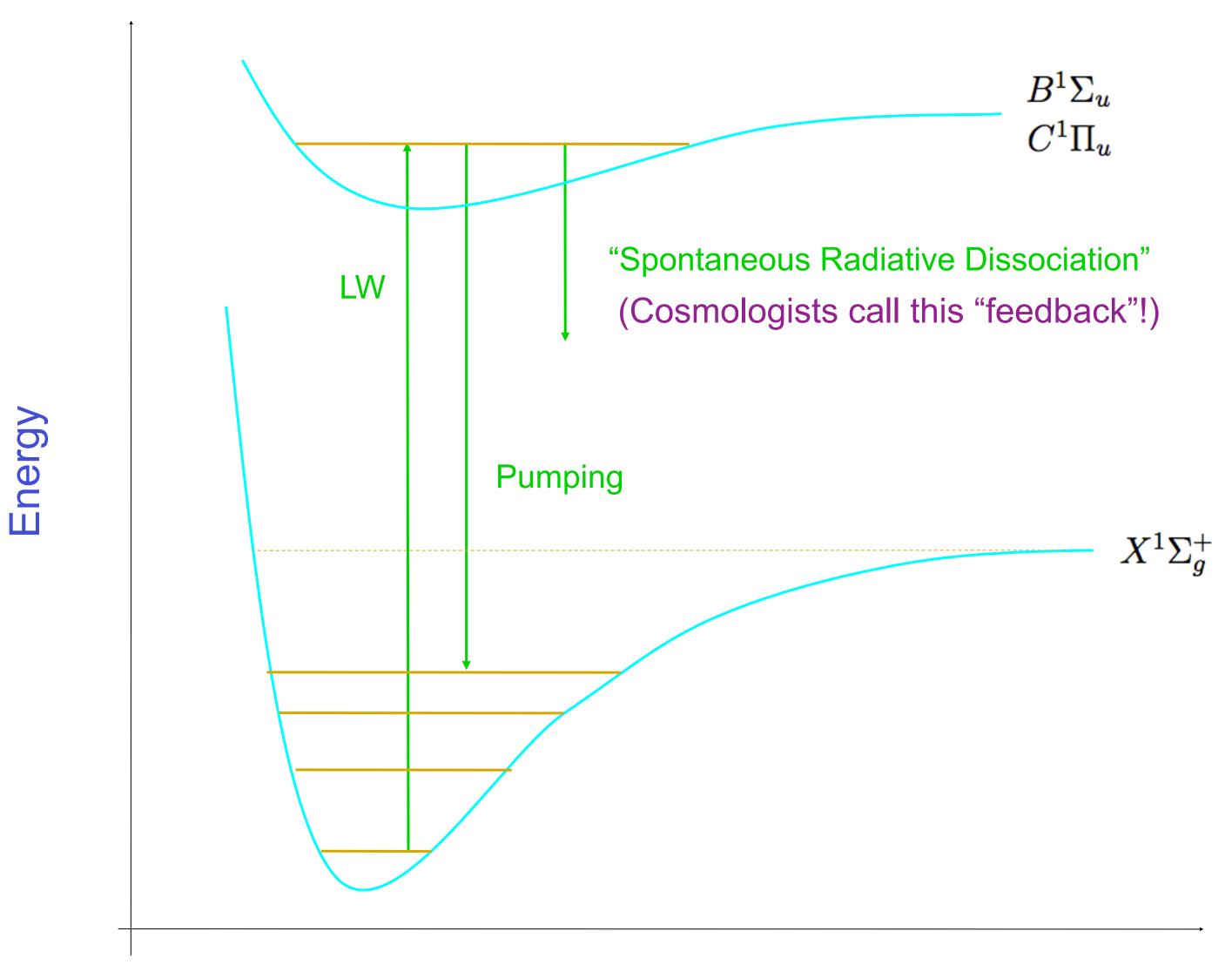


Internuclear Separation

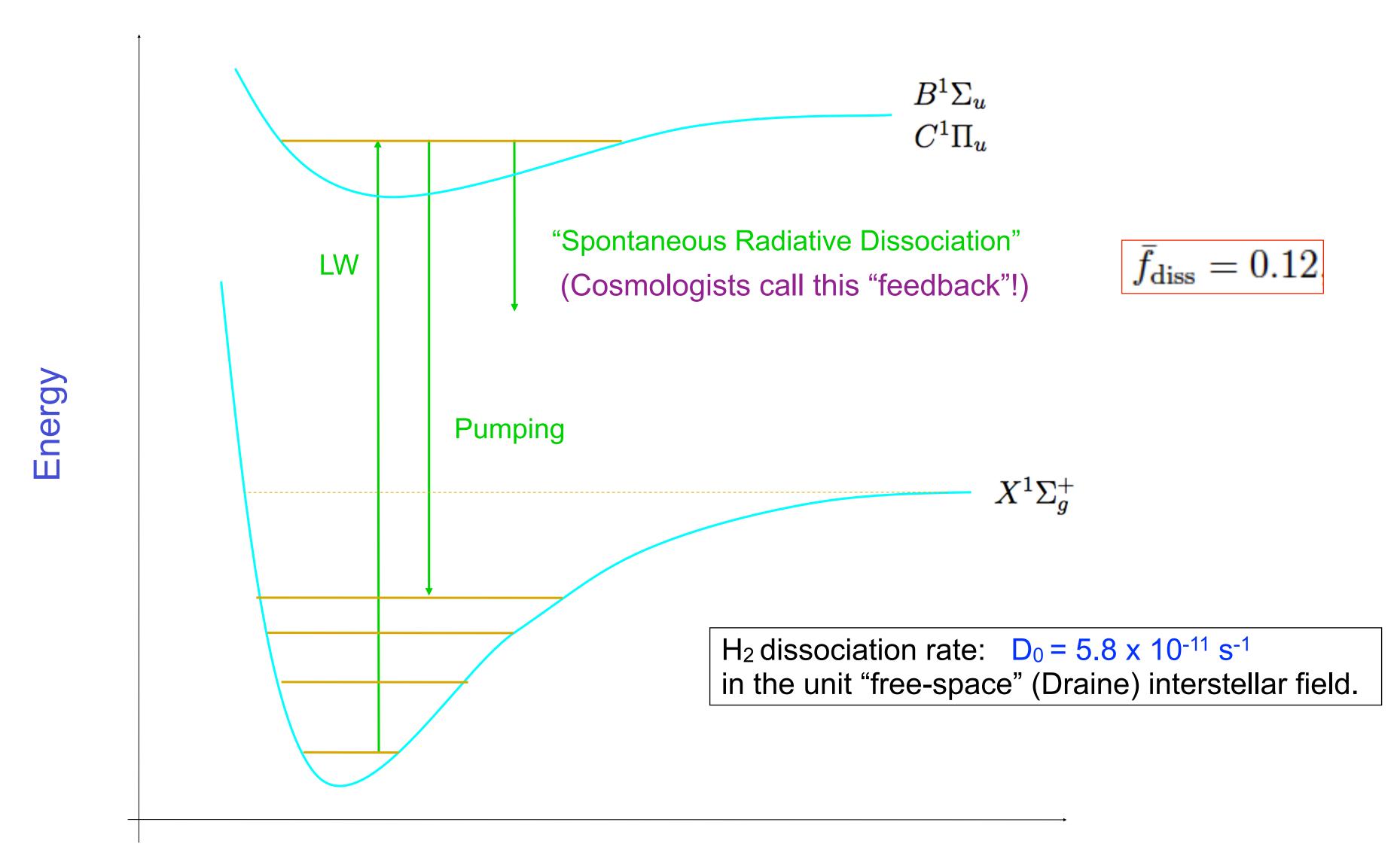


Internuclear Separation

Energy

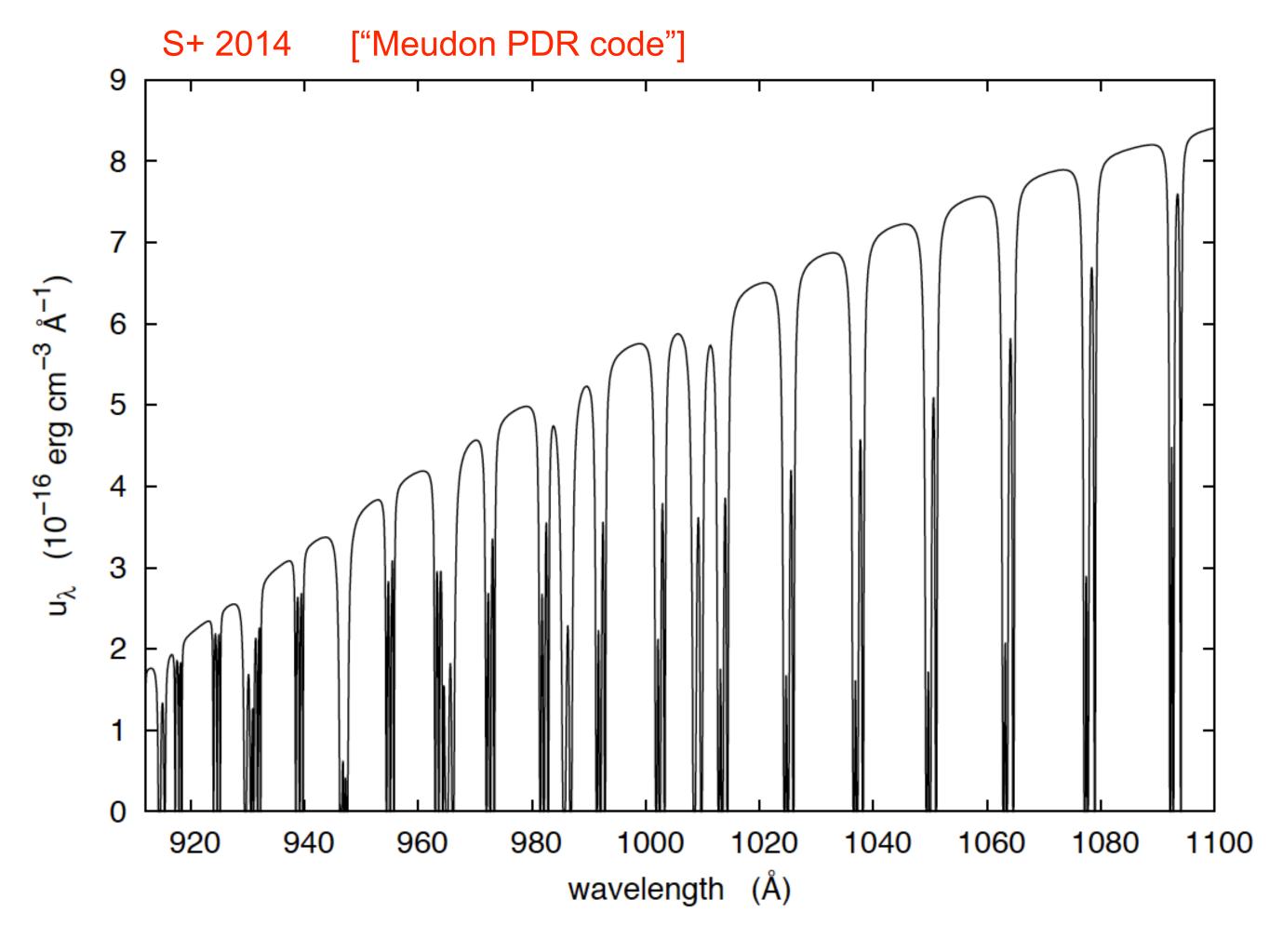


Internuclear Separation



Internuclear Separation

Lyman-Werner Radiative Transfer:



Characteristic multi-line H₂ absorption spectrum, and "self-shielding".

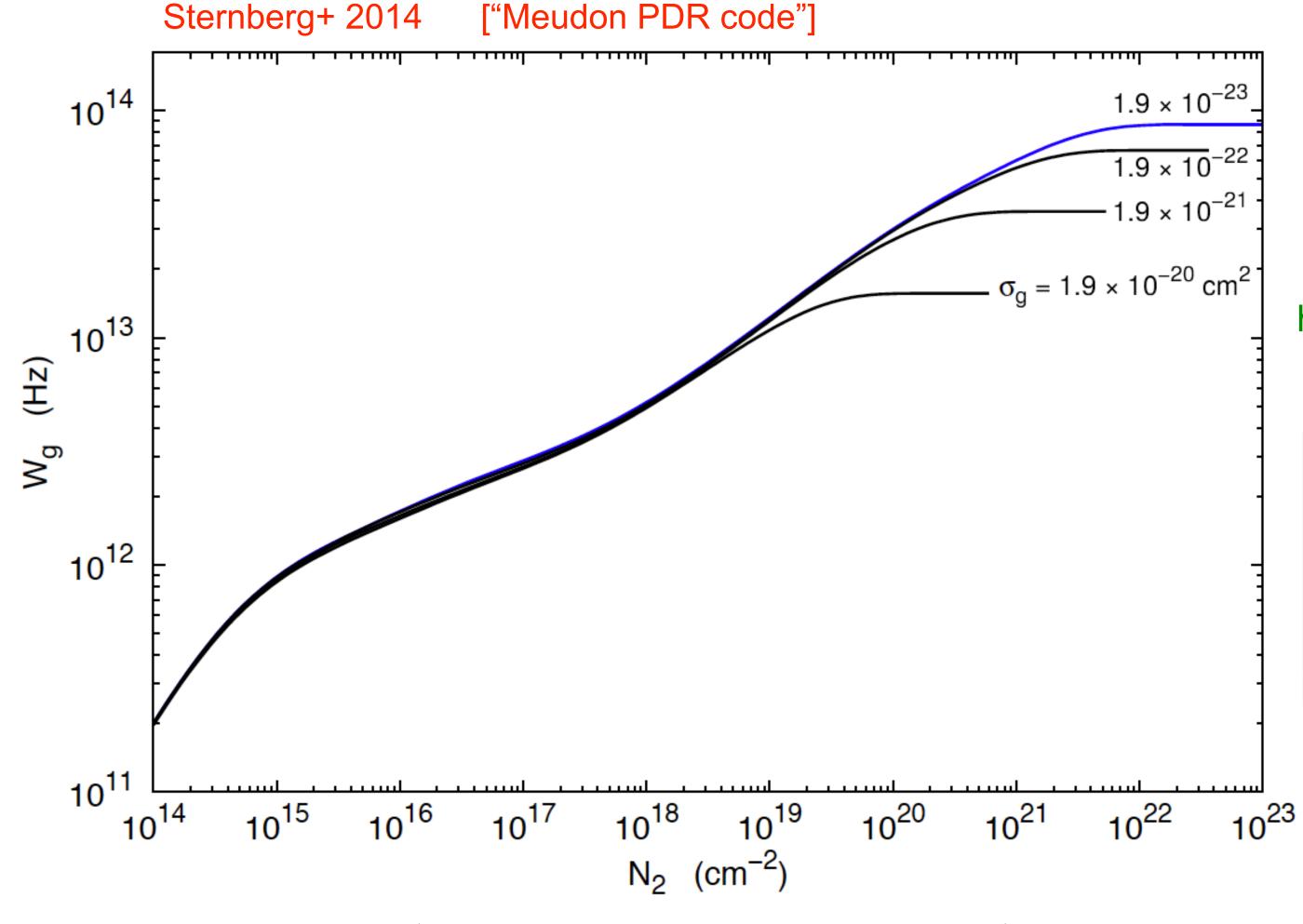
Dust absorption cross-section per hydrogen nucleus:

$$\sigma_g = 1.9 \times 10^{-21} \phi_g Z' \text{ cm}^2$$

Numerical radiative transfer on a fine frequency grid with a spectral resolution $\sim 10^5$.

Curve-of-Growth for the "H2 Dust-Limited Dissociation Bandwidth":

[integrated over all LW lines]



low metallicity

Z' "little w"

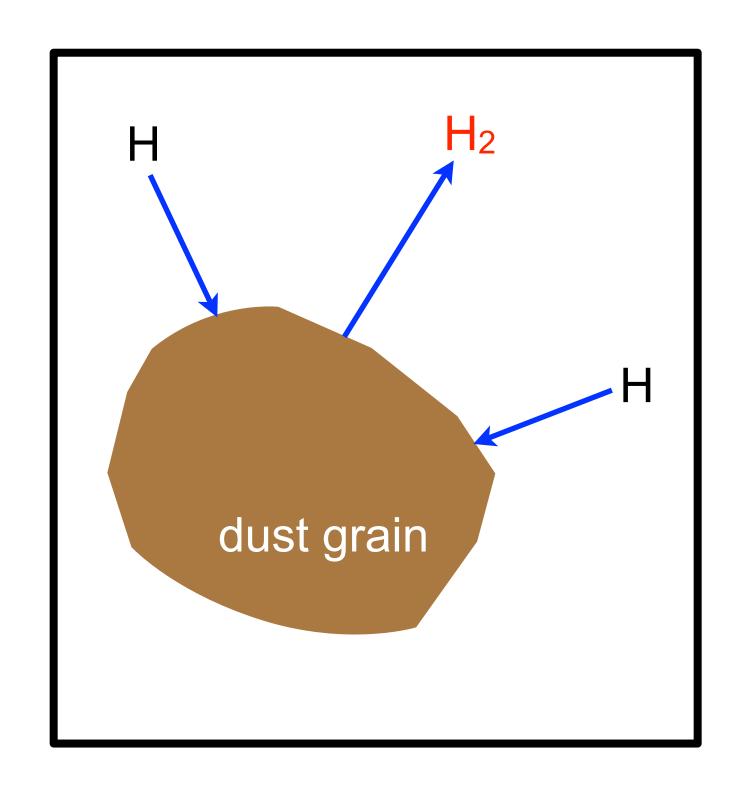
high metallicity

The bandwidth (Hz) of radiation absorbed in H₂ line dissociations in a dusty and fully molecular cloud.

"H2 dust" versus "H2 lines"

Universal: independent of radiation field intensity, or cloud gas density, etc.

H₂ Formation:



$$R = 3 \times 10^{-17} \left(\frac{T}{100 \text{ K}}\right)^{1/2} Z' \text{ cm}^3 \text{ s}^{-1}$$

In the absence of dust:

$$H + e \rightarrow H^- + photon$$

$$H^- + H \rightarrow H_2 + e$$

$$t_{\rm eq} = \frac{1}{Rn} = \frac{10^9}{Z'n} \, \mathrm{yr}$$

time scale for equilibrium

Basic Theory Question:

What is the total HI column density (cm⁻²) or HI mass surface density (M_☉ pc⁻²) in far-UV irradiated systems?

Dimensionless Parameter:

Sternberg 1988; McKee & Krumholz 2010; Sternberg+ 2014

$$\alpha G \equiv \bar{f}_{\rm diss} \frac{\sigma_g w F_0}{Rn}$$

Physical Meaning:

self-shielded H₂ dissociation rate

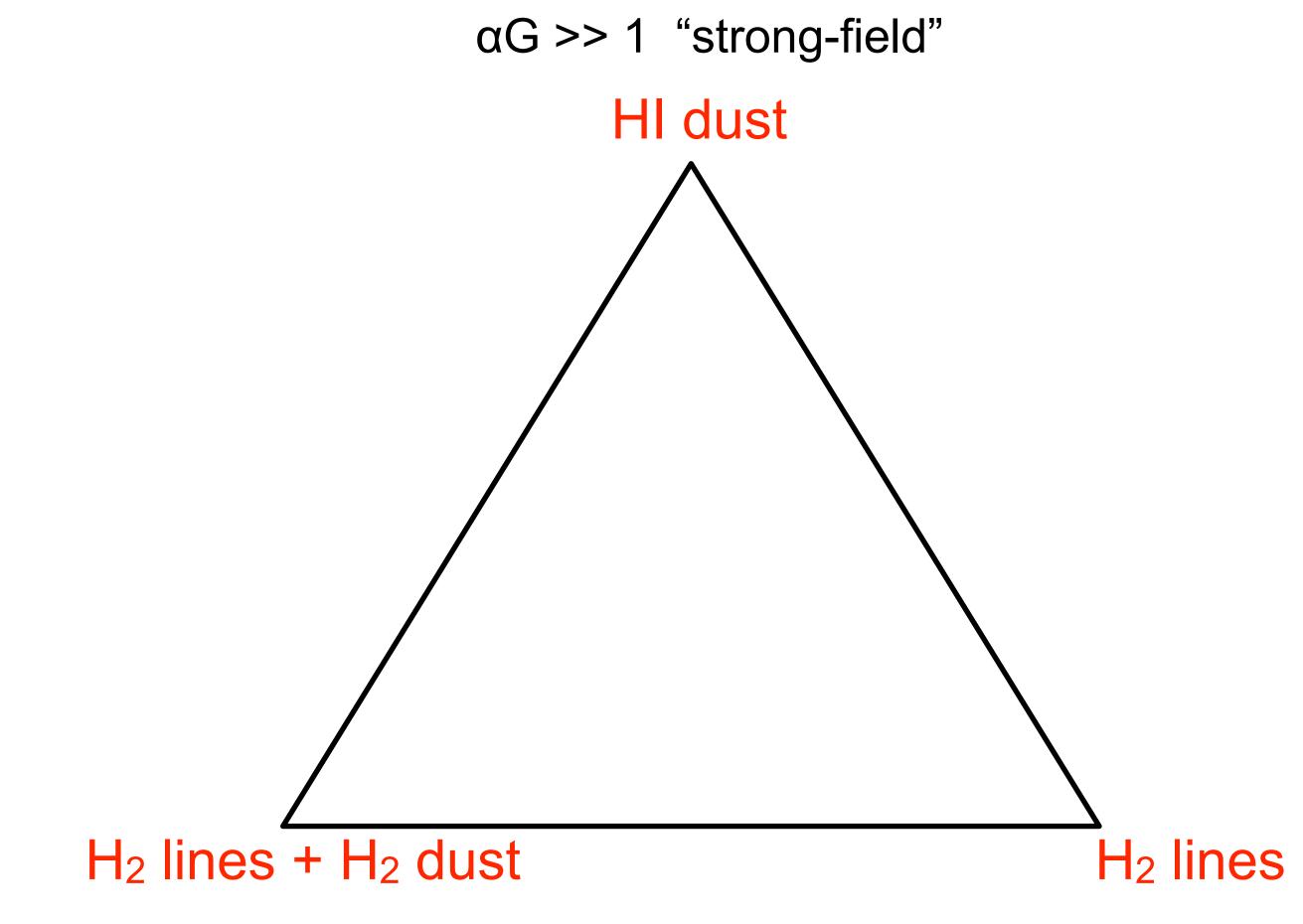
H₂ rate formation rate

or:

HI-dust absorption rate of the effective dissociation flux

free space H₂ photodissociation rate

Three-Way Competition for the FUV Absorption:



αG << 1 "weak-field" Z' ≥ 1 "high-metallicity"

αG << 1 "weak-field" Z' ≤ 1 "low-metallicity"

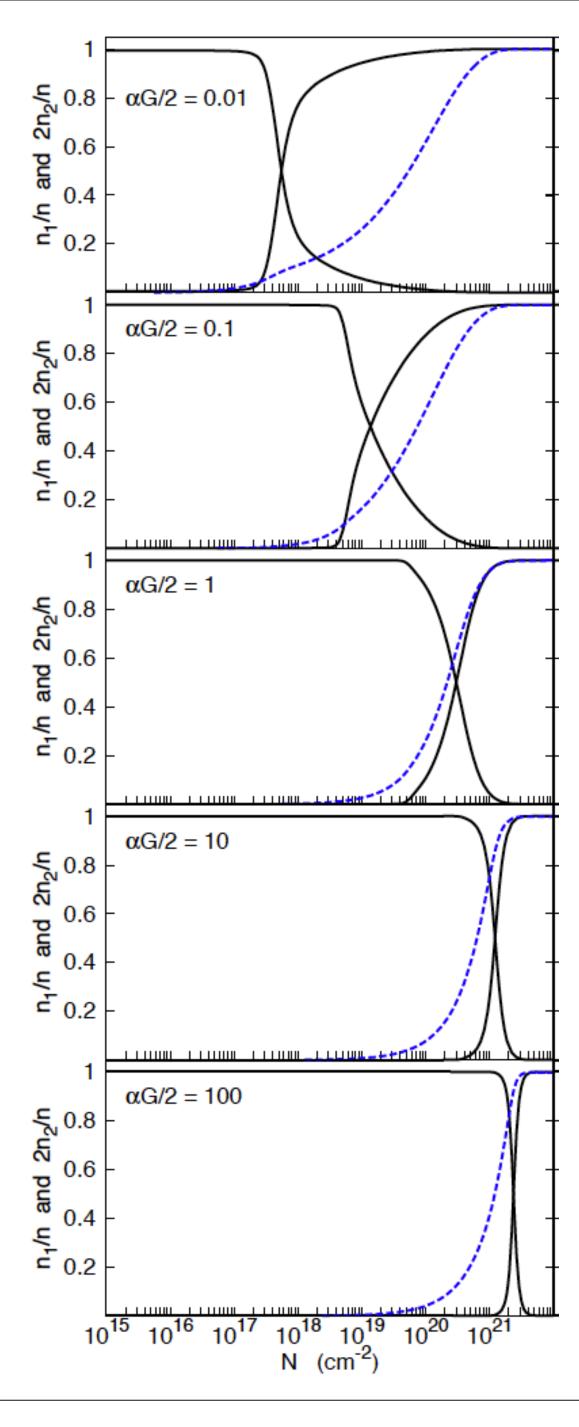
HI-to-H₂ Transition Profiles: Z'=1 (solar metallicity)

weak-field: most of the HI built up inside the H₂ zone, "gradual profiles"

FUV attenuation dominated by H₂ lines plus "H₂ dust"

strong-field: most of the HI built up in an outer layer, "sharp profiles"

FUV attenuation dominated by "HI-dust"



HI-to-H₂ Transition Profiles: Z'=1 (solar metallicity)

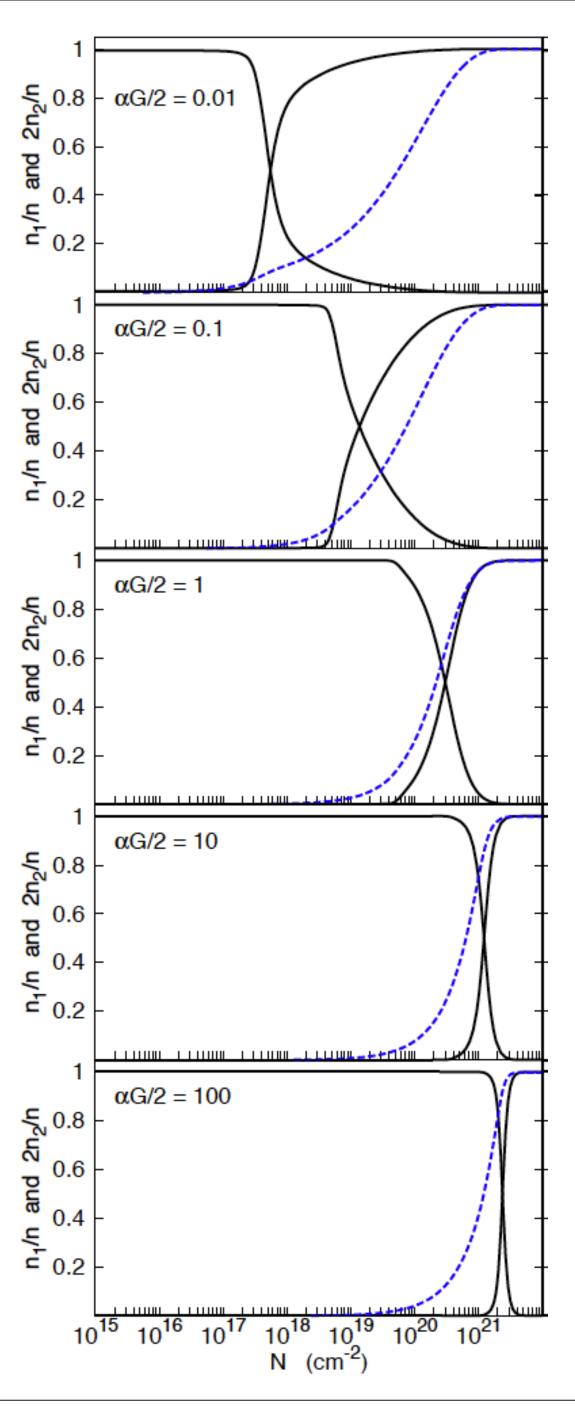
weak-field: most of the HI built up inside the H₂ zone, "gradual profiles"

FUV attenuation dominated by H₂ lines plus "H₂ dust"

For a given Z', the profile shapes depend on the dimensionless parameter αG .

strong-field: most of the HI built up in an outer layer, "sharp profiles"

FUV attenuation dominated by "HI-dust"



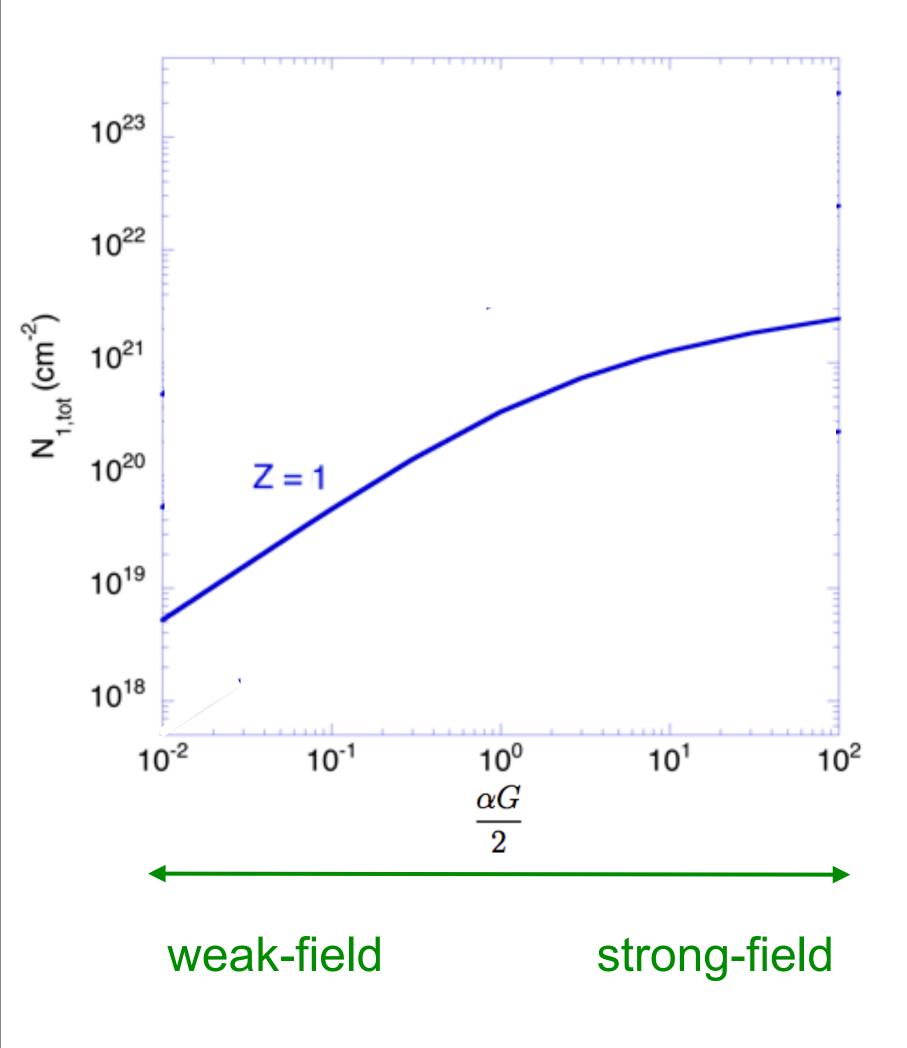
For a given αG the profile "scale lengths" are determined by Z'.

$$\sigma_g = 1.9 \times 10^{-21} \phi_g \ Z' \quad \text{cm}^2$$
 E.g., HI profiles for $\alpha G/2 \approx 1$
$$0.8 \quad 0.6 \quad Z' = 10 \quad 1 \quad 0.1 \quad 0.01$$

$$0.4 \quad 0.2 \quad 0.9 \quad 0.9$$

General Purpose Analytic Formula for the Total HI Column Density:

$$N_{1,\text{tot}} = \frac{1}{\sigma_g} \ln \left[\frac{\alpha G}{4} + 1 \right] = \frac{1}{\sigma_g} \ln \left[\frac{1}{4} \frac{\bar{f}_{\text{diss}} \sigma_g w F_0}{Rn} + 1 \right]$$

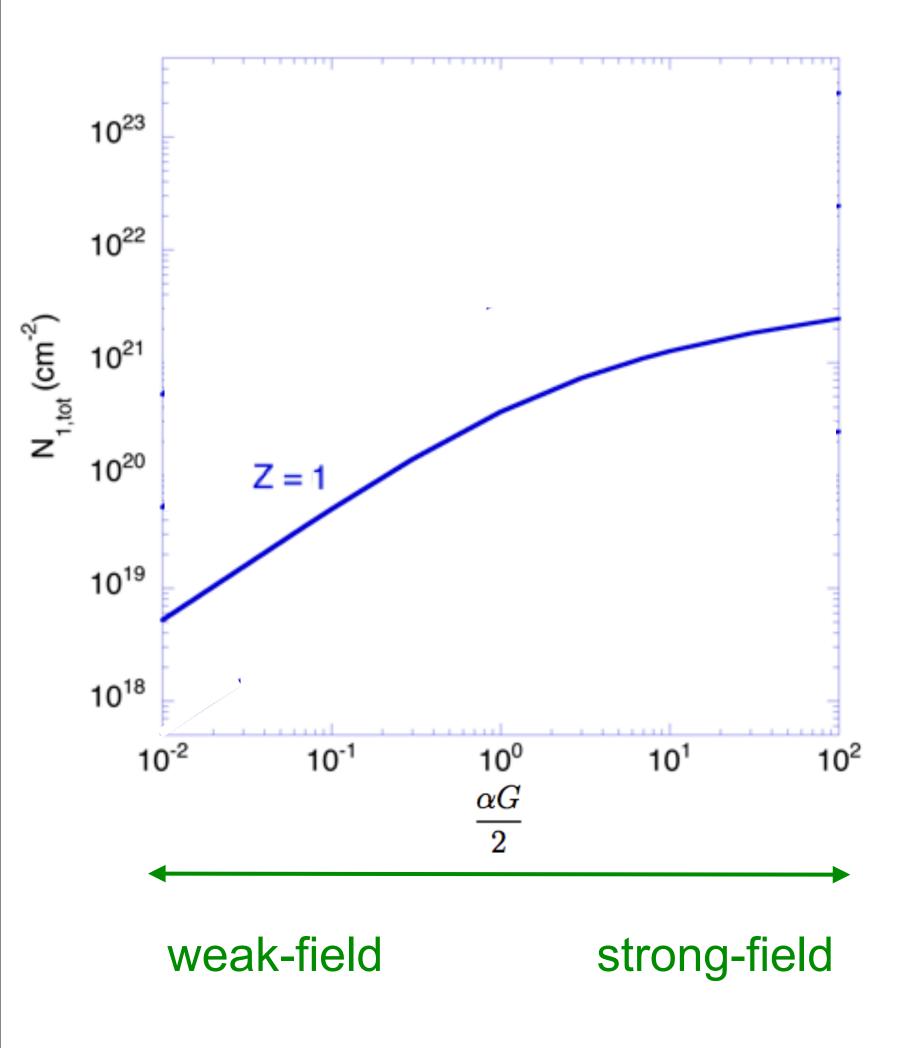


Valid for all regimes:

- weak and strong fields
- gradual to sharp transitions
- arbitrary metallicity

[note: no reference to the H₂ line photodissociation cross sections!]

Weak-Field Limit:



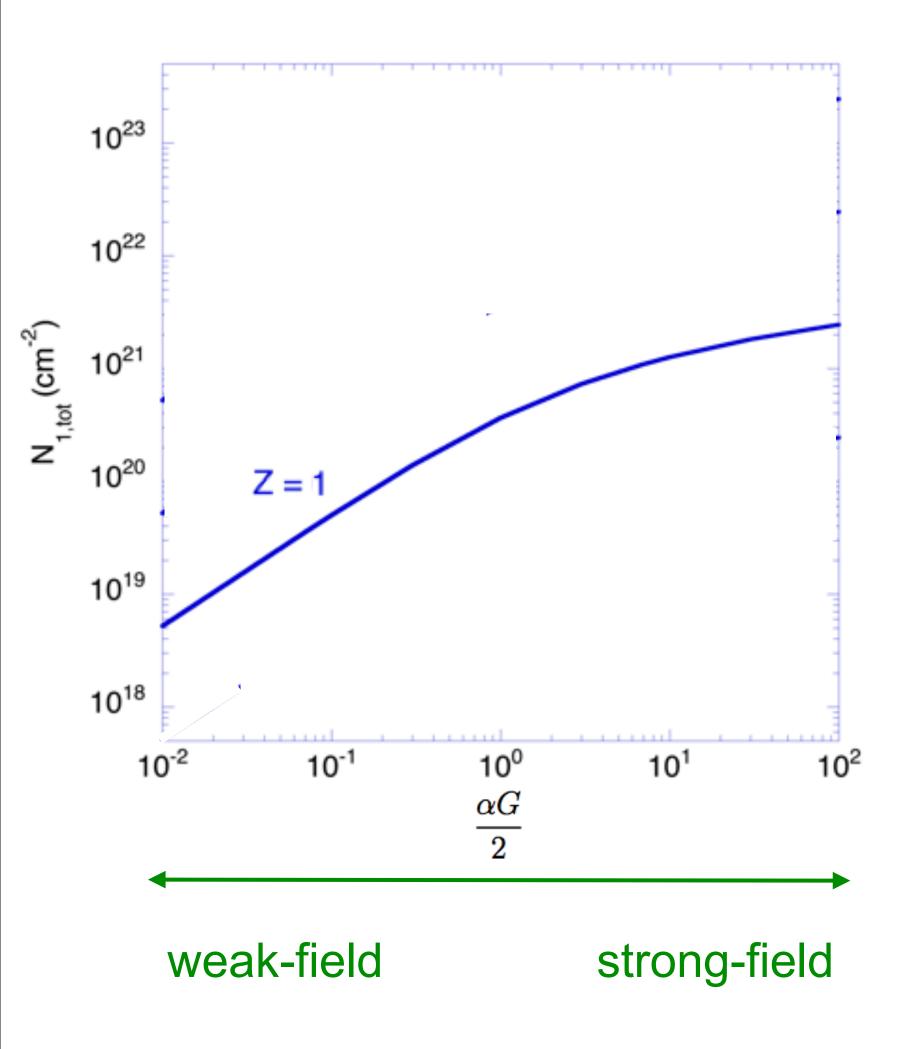
$$N_{1,\mathrm{tot}} = \frac{1}{4} \frac{\bar{f}_{\mathrm{diss}} w F_0}{Rn}$$

$$Rn N_{1, ext{tot}} = rac{1}{4} \bar{f}_{ ext{diss}} w F_0$$

formation rate per unit area = effective dissociation flux

(a "Strömgren Relation")

Strong-Field Limit:



$$N_{1,{
m tot}} pprox rac{1}{\sigma_g}$$
 (neglecting the logarithmic term)

makes sense: When HI-dust dominates the attenuation of the far-UV field, the HI-column is "self-limited" and

$$\tau_{\rm HIdust} = \sigma_g N_{1,\rm tot} \approx 1$$

Heavy-Element Abundances "Metallicity":

$$R \propto Z'$$
 $\sigma_g \propto Z'$

H₂ formation rate coefficient and dust absorption cross-section both proportional to metallicity.

$$N_{1,\mathrm{tot}} = \frac{1}{4} \frac{\bar{f}_{\mathrm{diss}} w F_0}{Rn} \propto \frac{1}{Z'}$$

 $N_{1,{
m tot}} \approx \frac{1}{\sigma_g} \propto \frac{1}{Z'}$

weak-field

strong-field

Heavy-Element Abundances "Metallicity":



$$R \propto Z'$$
 $\sigma_g \propto Z'$

H₂ formation rate coefficient and dust absorption cross-section both proportional to metallicity.

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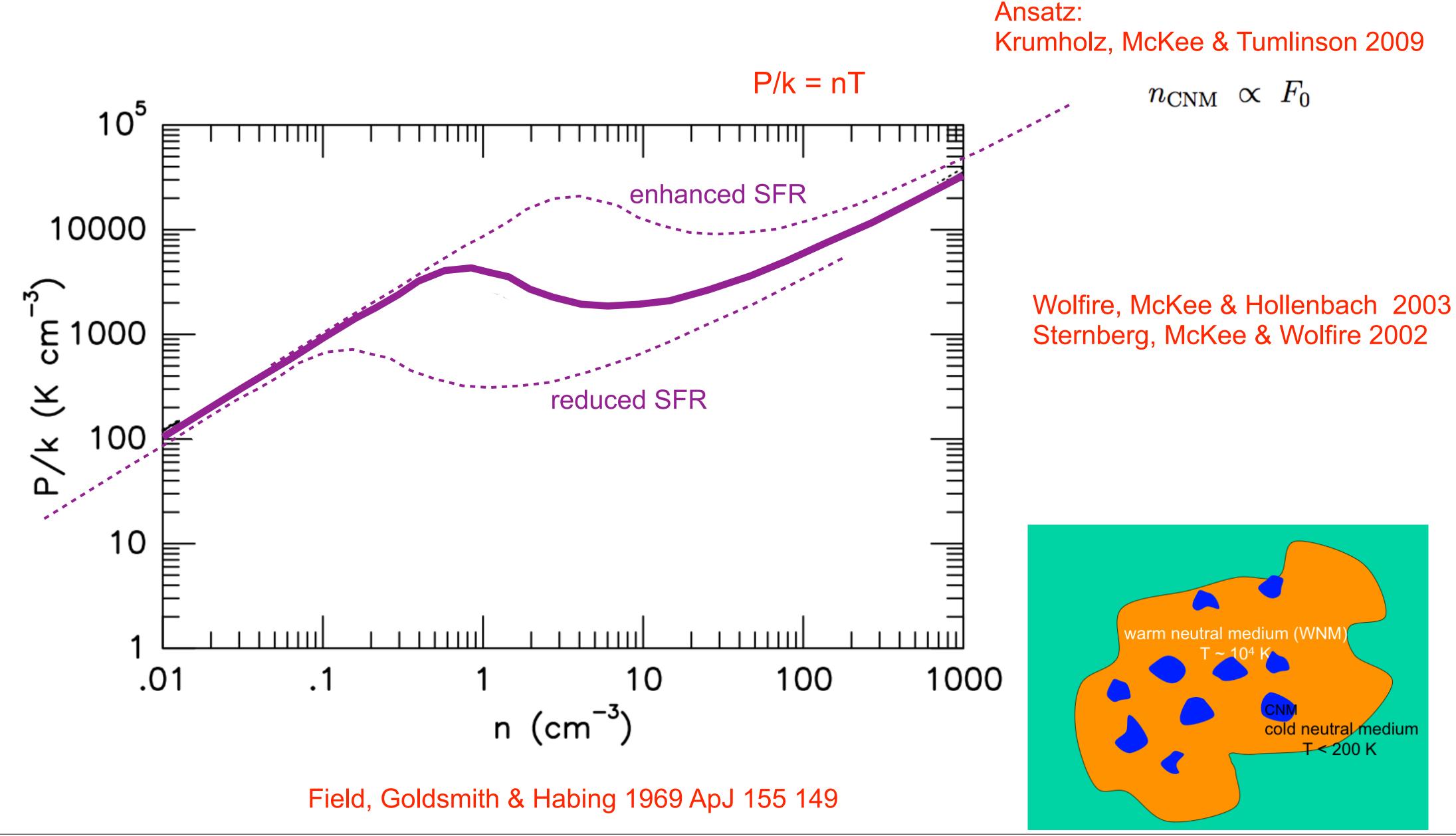
weak-field

strong-field

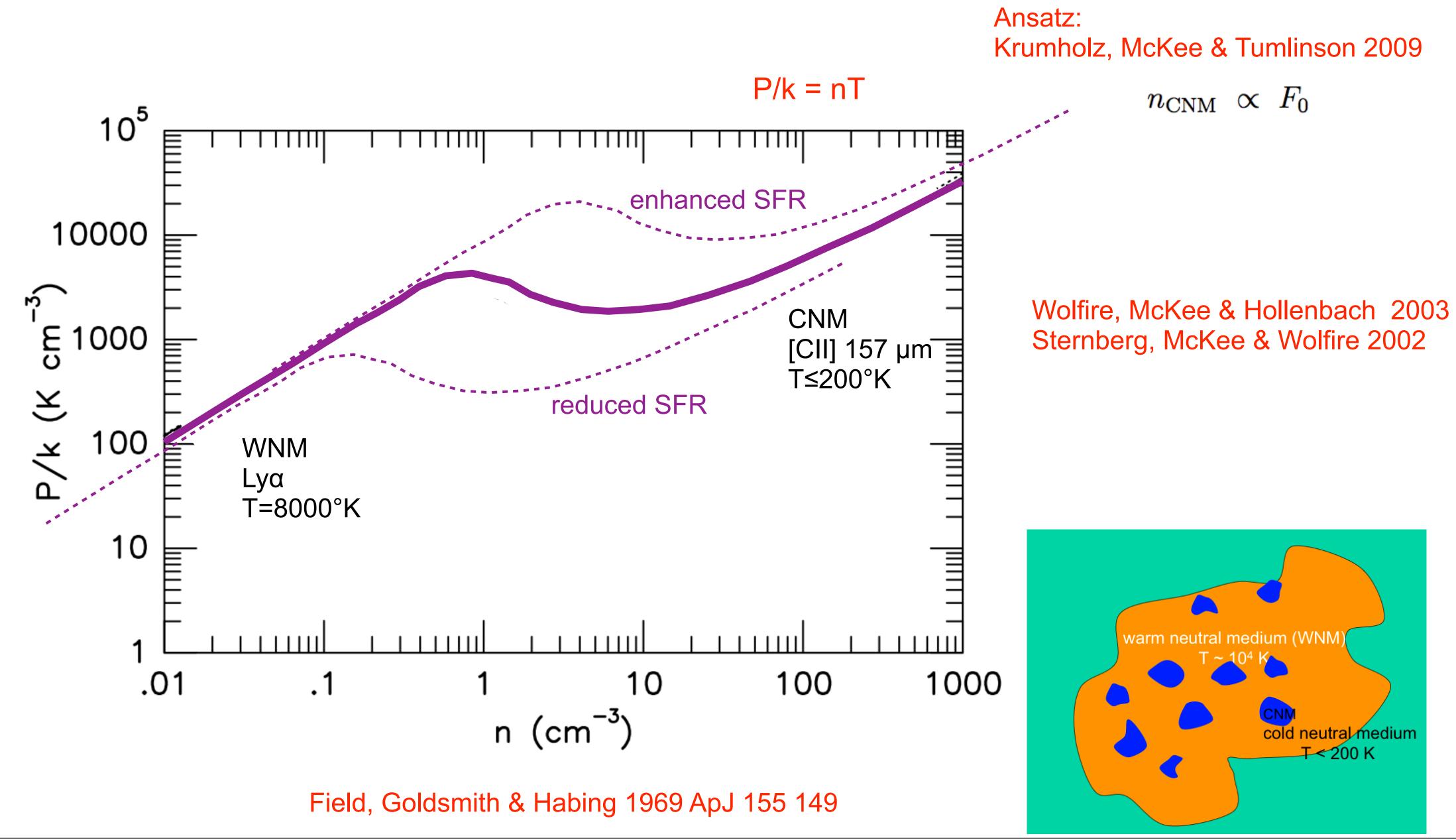
$$N_{1,\mathrm{tot}} = \frac{1}{\sigma_g} \ln \left[\frac{\alpha G}{4} + 1 \right] = \frac{1}{\sigma_g} \ln \left[\frac{1}{4} \frac{\bar{f}_{\mathrm{diss}} \sigma_g w F_0}{Rn} + 1 \right]$$

- useful for interpreting 21cm observations
- incorporation into hydrodynamics simulations
- application to "self-regulated" media

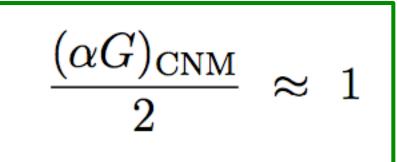
HI Thermal Phases in Self-Regulated Media:



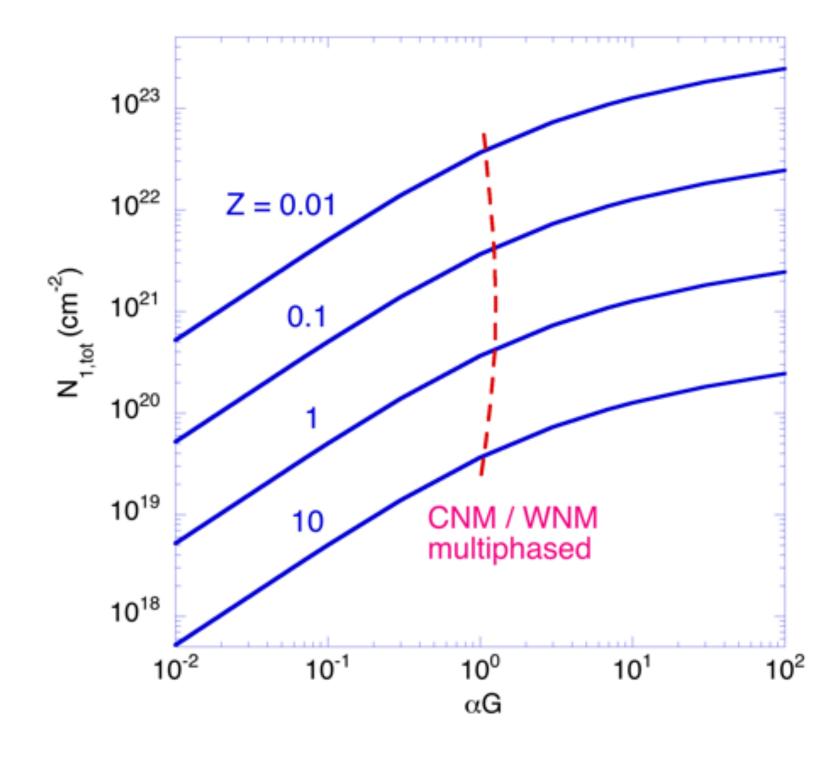
HI Thermal Phases in Self-Regulated Media:



HI Column Density for Self-Regulated Media:



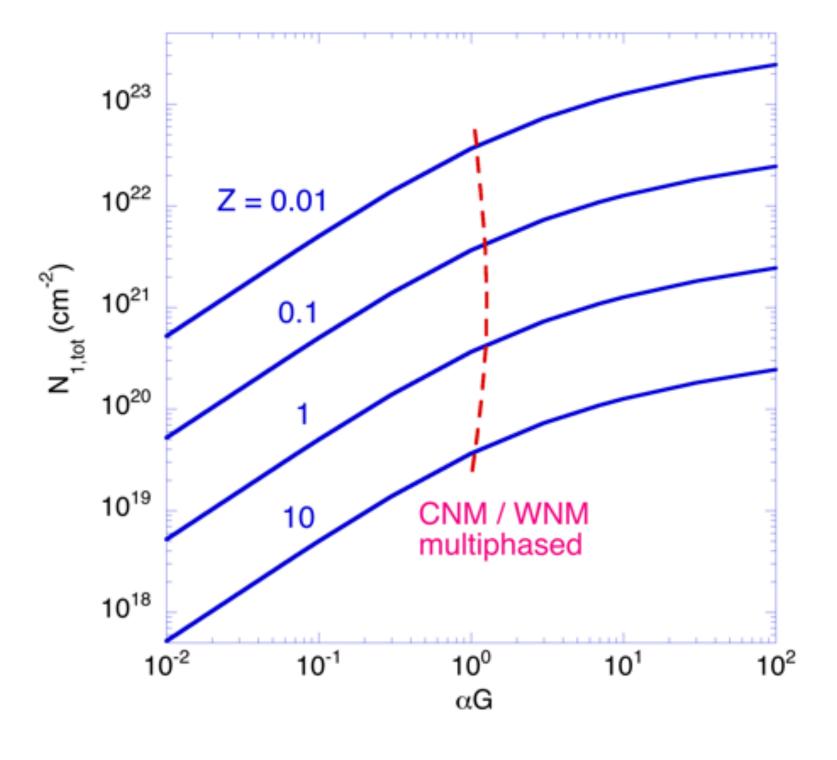
remarkable!



HI Column Density for Self-Regulated Media:

$$\frac{(\alpha G)_{\mathrm{CNM}}}{2} \approx 1$$

remarkable!



$$\Sigma_{\mathrm{HI}} \approx \frac{6}{\phi_g Z'} \quad \mathrm{M}_{\odot} \; \mathrm{pc}^{-2}$$

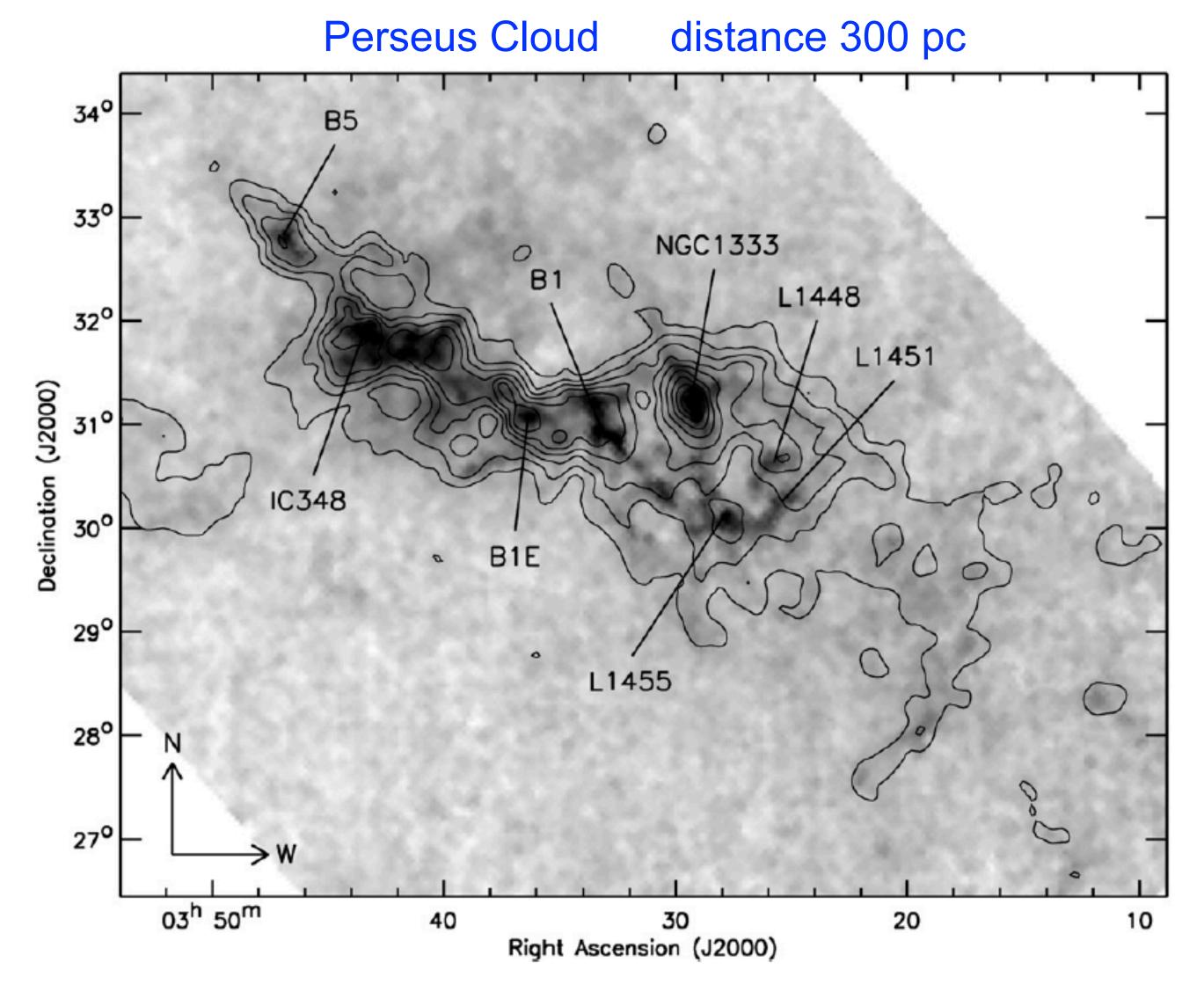
characteristic HI photodissociation mass surface density in self-regulated systems

$$\Sigma_{\rm gas,*}(Z') \equiv 2 \times \Sigma_{\rm HI} \approx \frac{12}{\phi_g Z'} \quad {\rm M}_{\odot} \ {\rm pc}^{-2}$$

"transition" total gas mass surface density... star-formation threshold...galaxies.

independent of radiation field intensity and/or gas density.

HI-to-H₂ in Perseus:



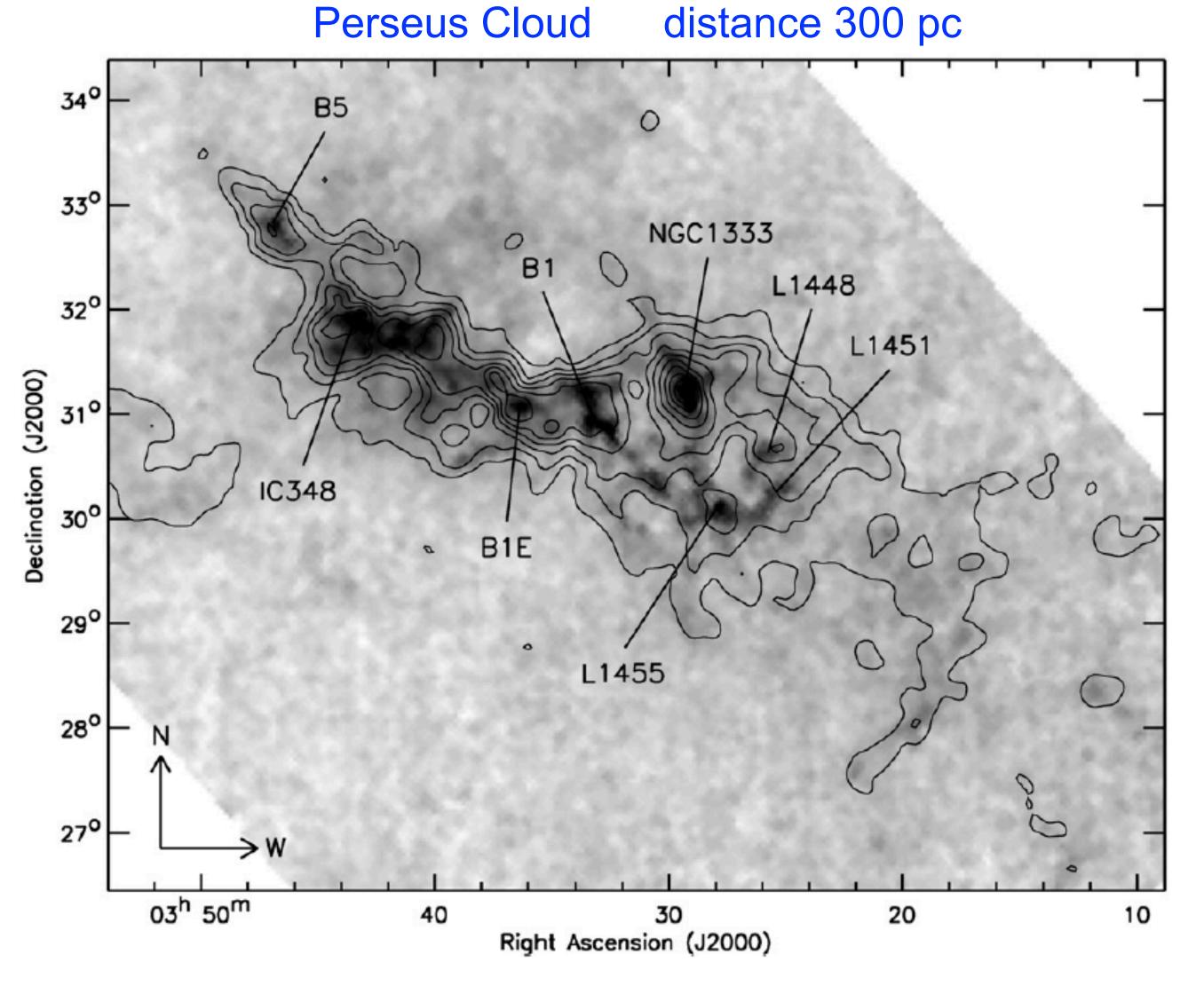
2MASS A_V and I_{CO} CfA & COMPLETE surveys

Dame+ 2001 Ridge+ 2006

HI-to-H₂ in Perseus:

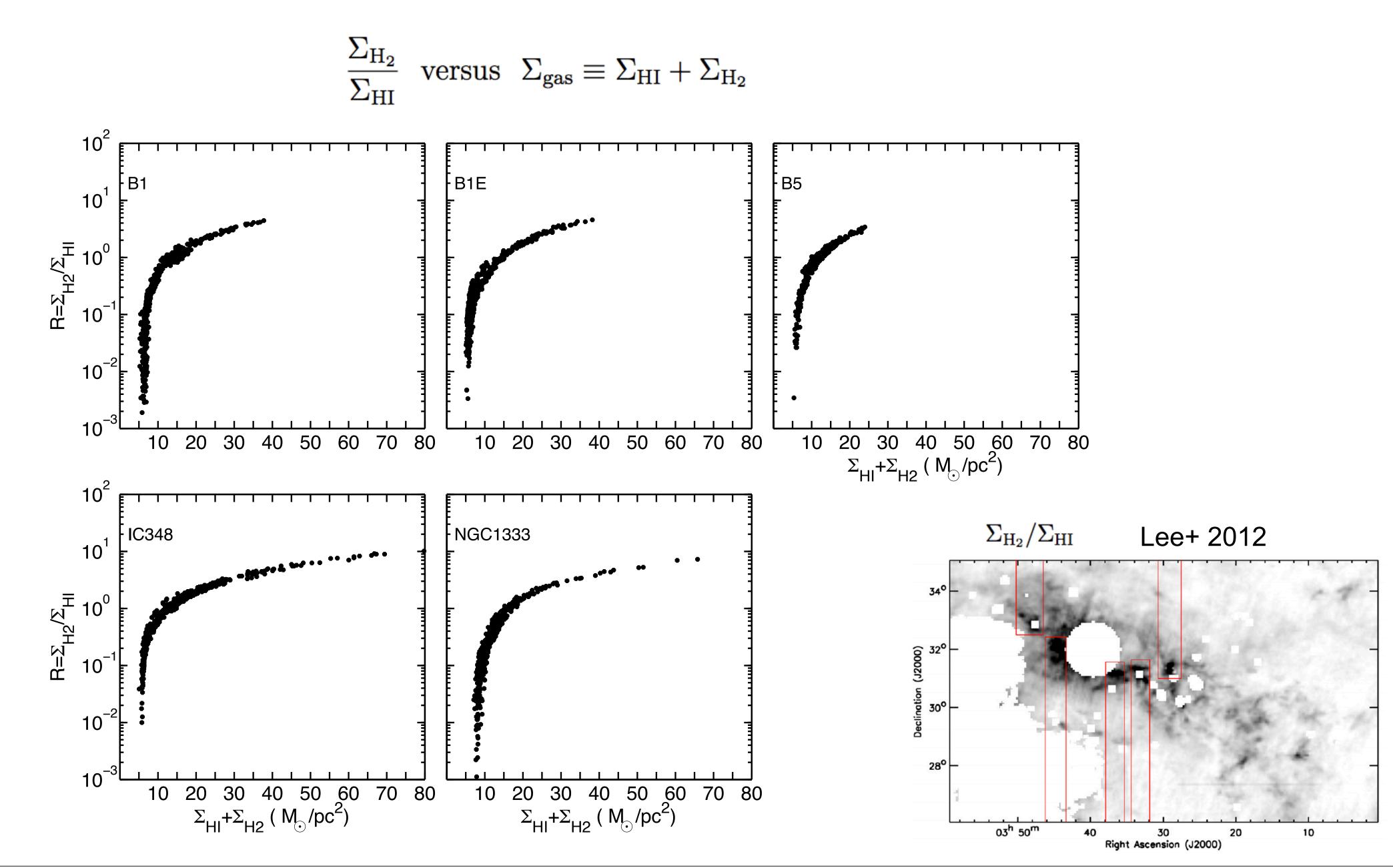
dark clouds & low-mass star-forming regions:

B1 B1E B5 NGC1333 IC348

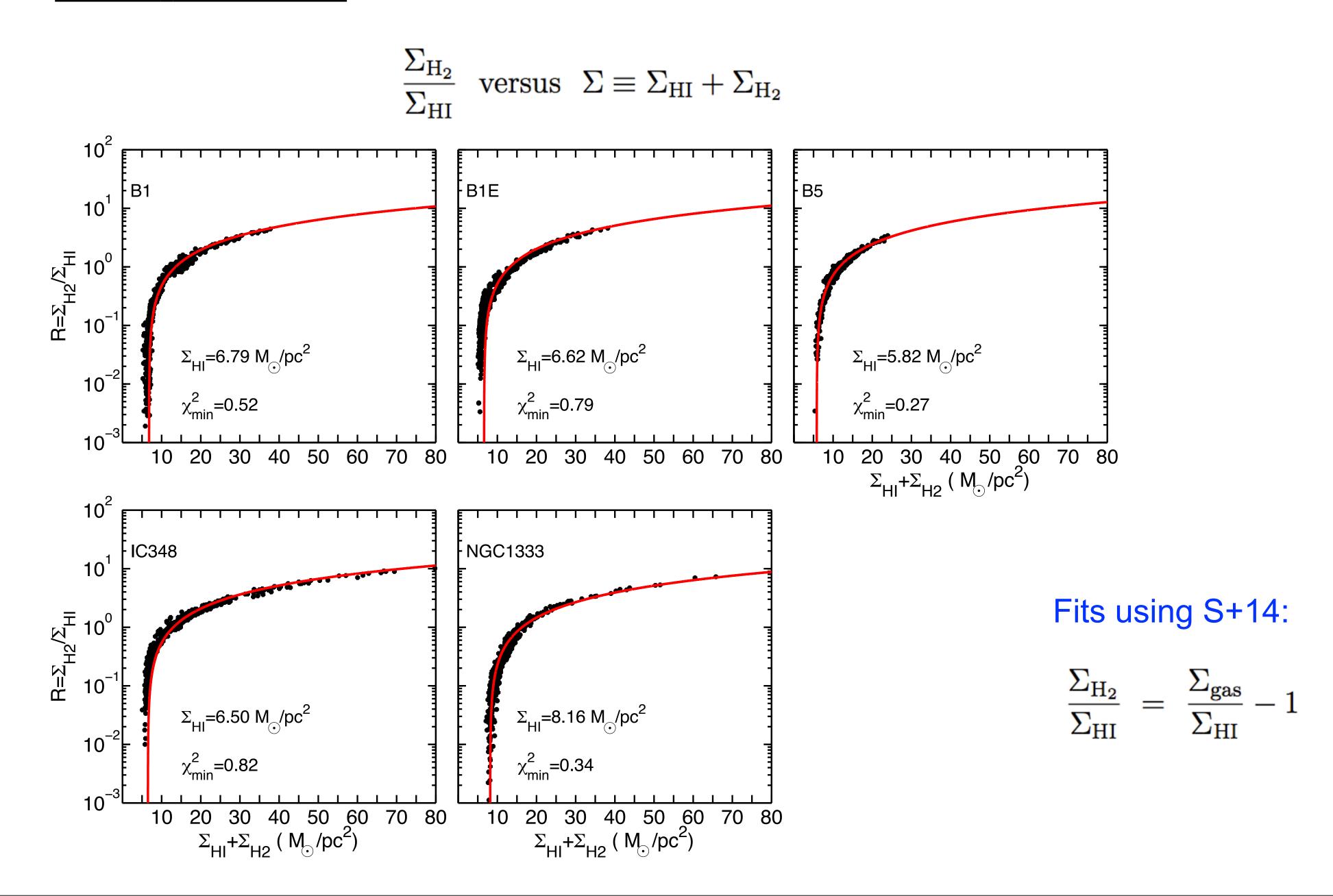


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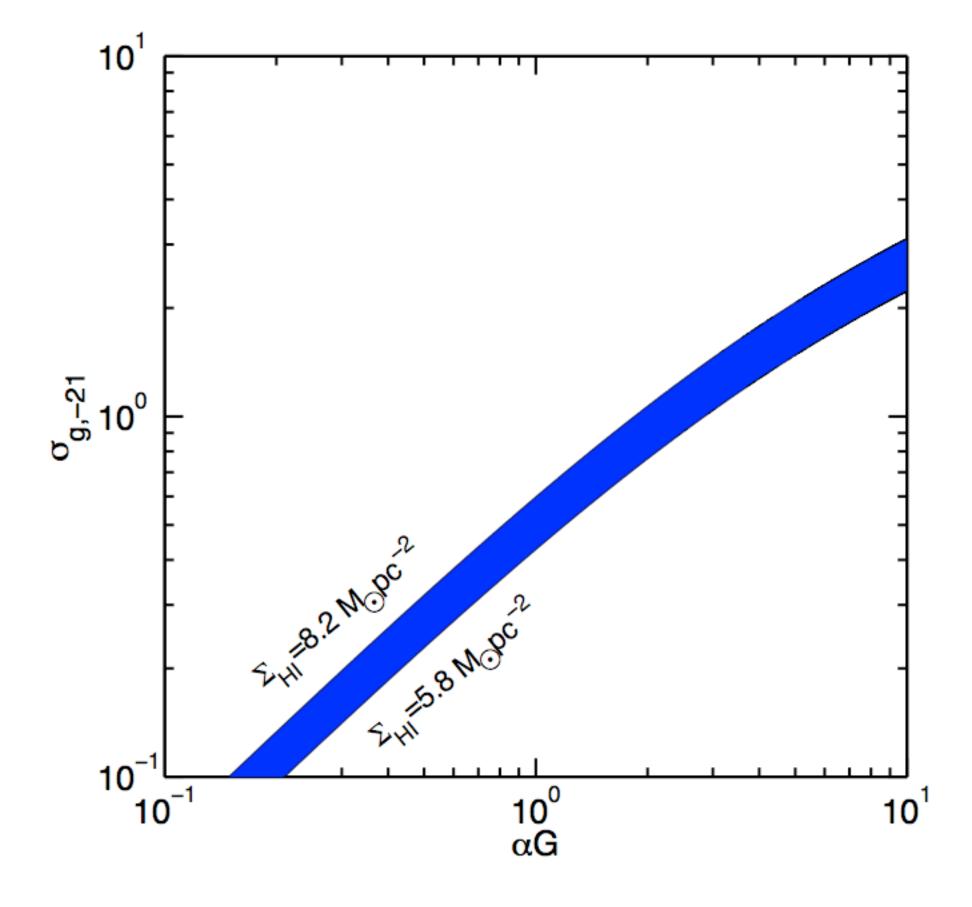
Dame+ 2001 Ridge+ 2006

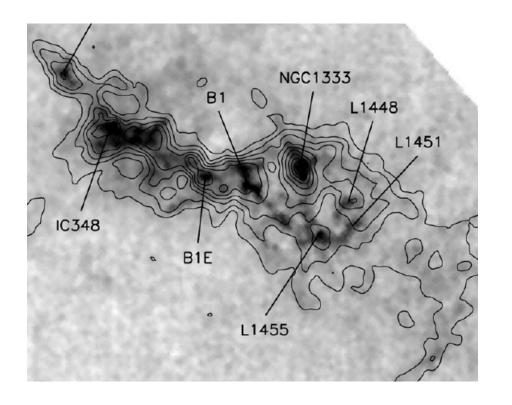


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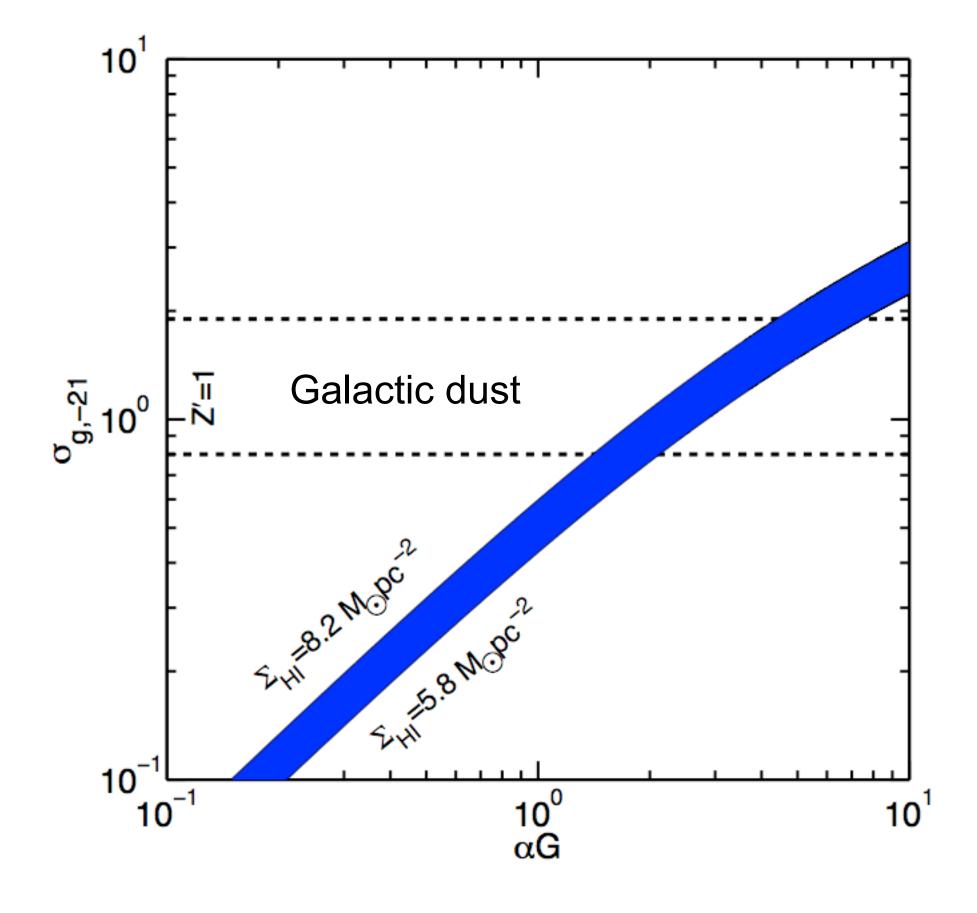


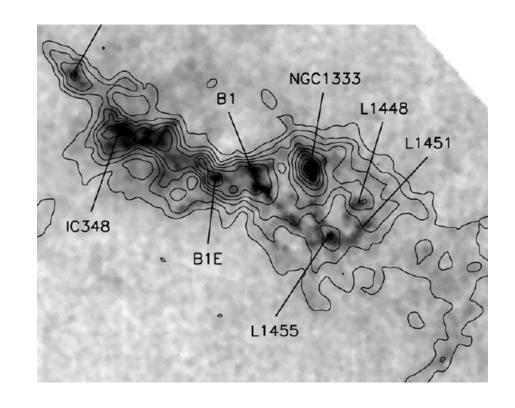
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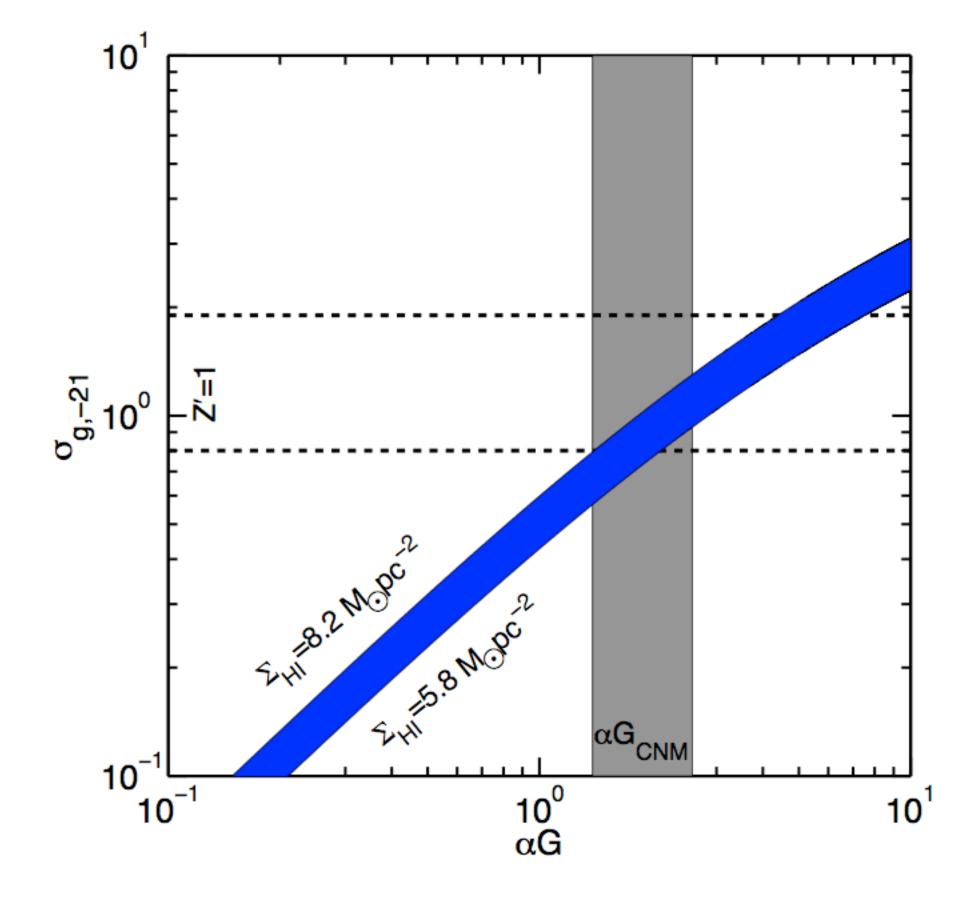


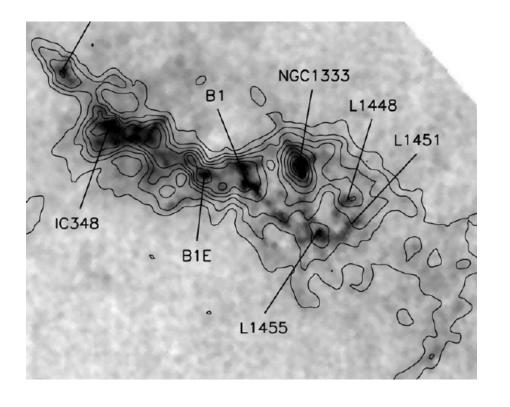
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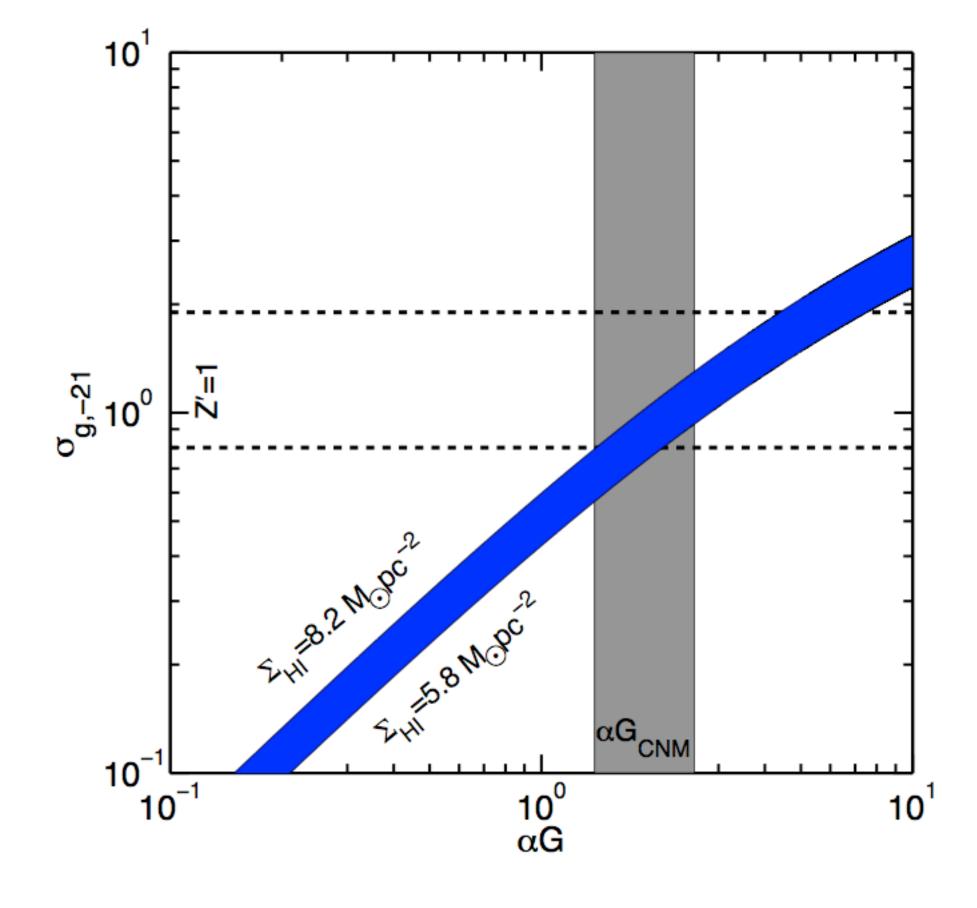


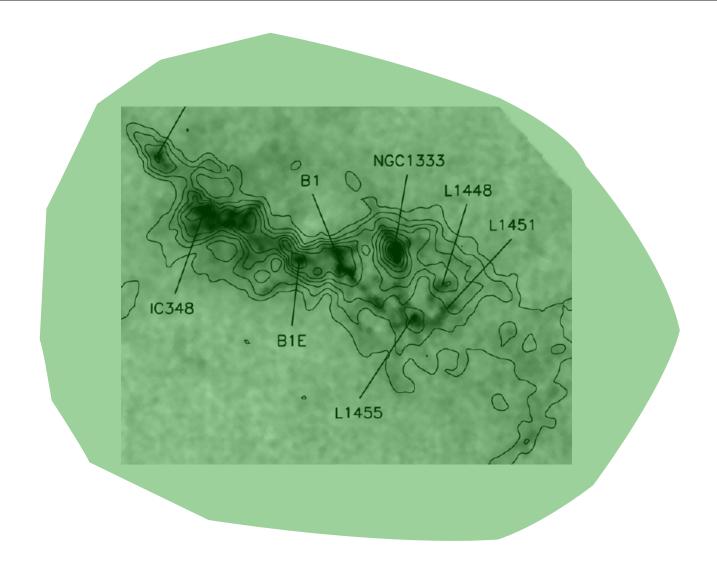
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$$N_{1,\mathrm{tot}} = \frac{1}{\sigma_g} \ln \left[\frac{\alpha G}{4} + 1 \right]$$





Galactic interstellar LW radiation field: $F_0 \approx 2 \times 10^7$ photons cm⁻² s⁻¹

Conclusions: $\alpha G \approx 2 \text{ to } 6$

- FUV absorption dominated by HI-dust
- $n_{HI} \approx 20 \text{ to } 3 \text{ cm}^{-3}$
- consistent with CNM/WNM mixture

Galaxies:

e.g. Bigiel+ 08

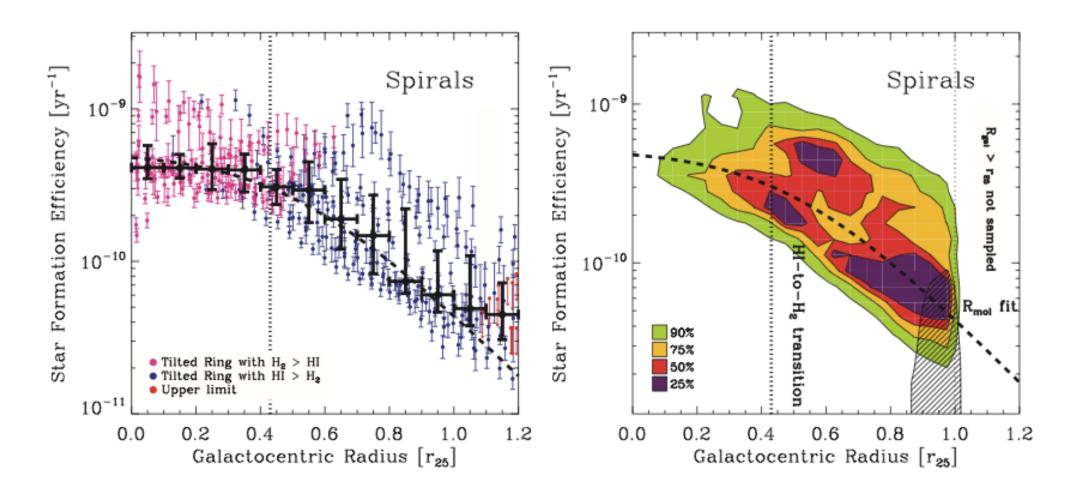
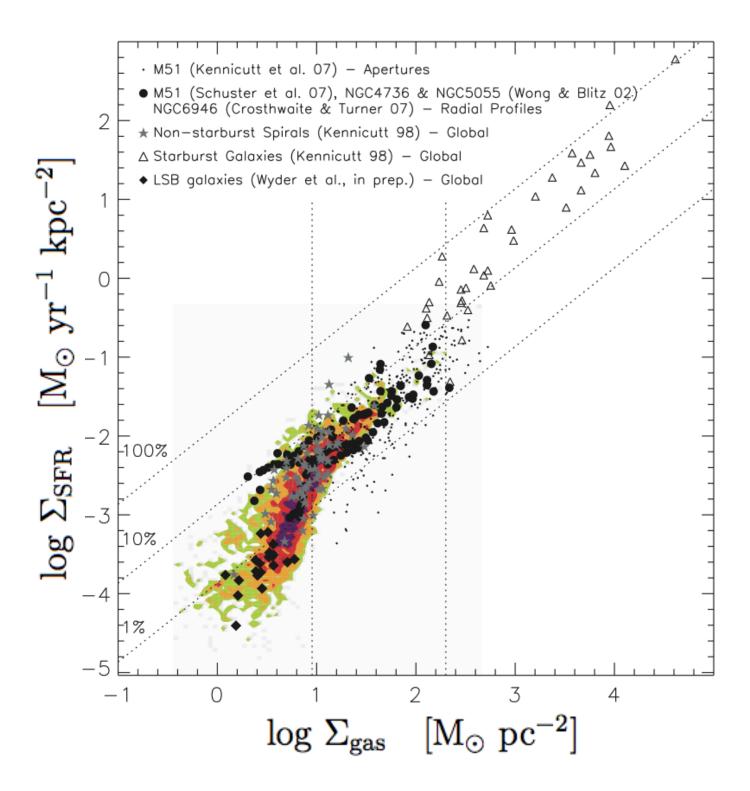


Table 5
Conditions at the H_I-to-H₂ Transition

Quantity	Median Value ^a	Scatter	Scatter in log ₁₀
$r_{\rm gal} (r_{25})$	0.43	0.18	0.17
$\Sigma_* (M_{\odot} \text{ pc}^{-2})$	81	25	0.15
$\Sigma_{\rm gas}~(M_{\odot}~{ m pc}^{-2})$	14	6	0.18
$P_h/k_{\rm B}~({\rm cm}^{-3}~{\rm K})$	2.3×10^{4}	1.5×10^{4}	0.26
$\tau_{\rm ff}$ (yr)	4.2×10^{7}	1.2×10^{7}	0.14
$\tau_{\rm orb}$ (yr)	1.8×10^{8}	0.4×10^{8}	0.09
$Q_{ m gas}$	3.8	2.6	0.31
$Q_{ m stars+gas}$	1.6	0.4	0.09



Assuming
$$\alpha G = (\alpha G)_{\rm CNM}$$

$$\Sigma_{\rm gas,*} \approx \frac{12}{\phi_g Z'} \quad {\rm M}_{\odot} \; {\rm pc}^{-2}$$

Galaxies:

e.g. Bigiel+ 08

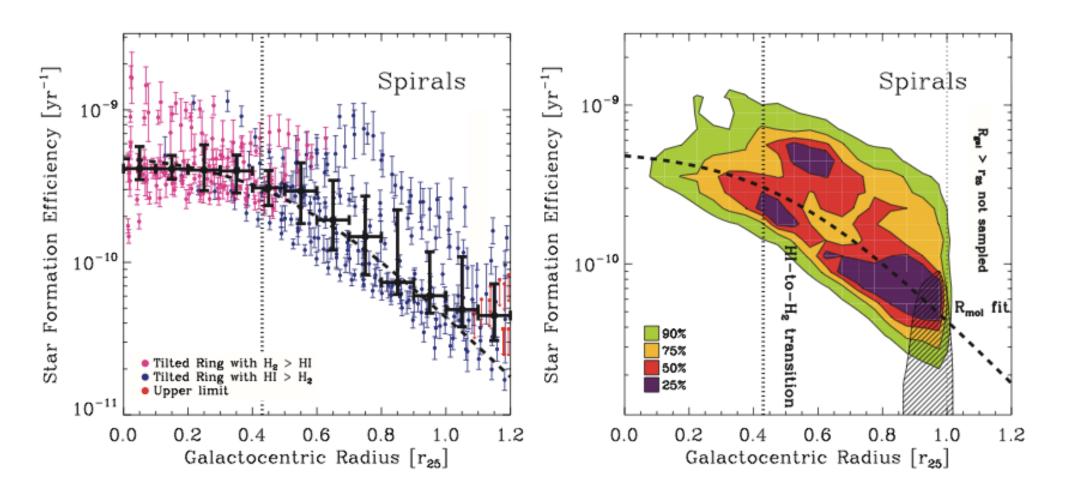
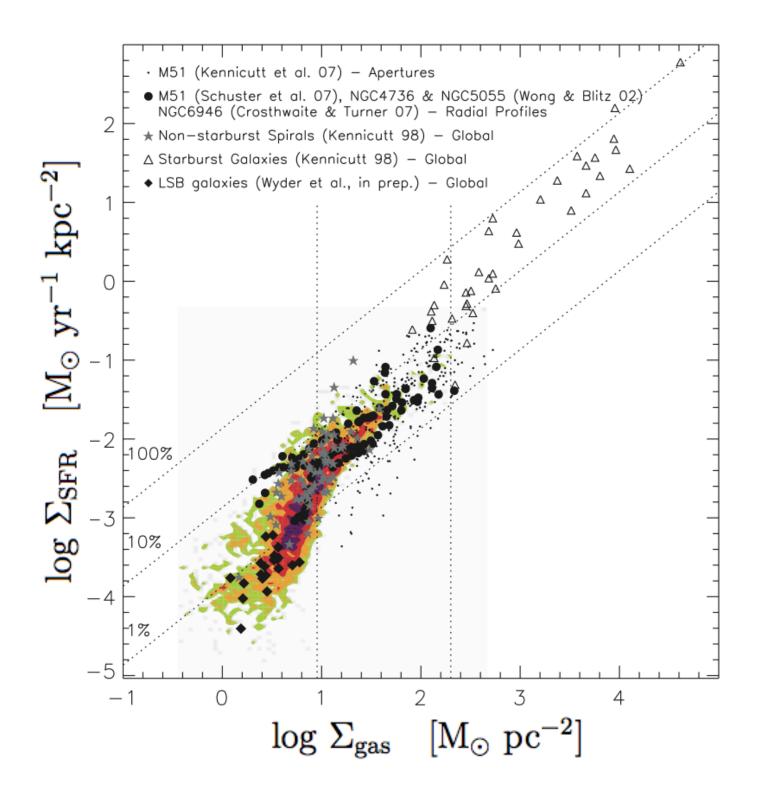


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Caveat: This interpretation requires typically "one" primary cloud per line-of-sight.

To Conclude:

$$N_{1,\text{tot}} = \frac{1}{\sigma_g} \ln \left[\frac{\alpha G}{4} + 1 \right] = \frac{1}{\sigma_g} \ln \left[\frac{1}{4} \frac{\bar{f}_{\text{diss}} \sigma_g w F_0}{Rn} + 1 \right]$$

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